# Galactic Dark Matter Signature in IceCube DeepCore

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**Cross Strait Meeting** 

F.-F. Lee and G.-L. Lin, Phys. Rev D85, 023529 (2012) F.-F. Lee G.-L Lin and Sming Tsai, in preparation



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# **Outline**

# Introduction

- IceCube 22-string search for galactic dark matter
- Detection sensitivity in IceCube+DeepCore phenomenological analysis
- Constraints from gamma ray observations
- ♦ Summary

# **Dark Matter**

An unknown type of matter does neither emit nor reflect EM radiation.

• Weakly Interacting Massive Particles (WIMPs) are one of the leading candidates for DM.

• WIMPs are theoretically well motivated and capable of producing the correct relic density.





# **IceCube Neutrino Observatory**



#### **Track Events & Cascade Events**

#### **Track Events**

Charged - Current  $\nu_{\mu}$  interaction :  $\nu_{\mu} + N \rightarrow \mu^{-} + X$ 

#### **Cascade Events**

Neutral - Current  $v_l$  interaction :  $v_l + N \rightarrow v_l + X$ (Hadronic)

Charged - Current  $\nu_{e}$  interacti**a** :  $\nu_{e} + N \rightarrow e^{-}(EM) + X(Hadronic)$ 

Charged - Current  $v_{\tau}$  interaction :  $v_{\tau} + N \rightarrow \tau^- + X$ 

# IceCube 22 String Result



# Phys. Rev. D 84, 022004 (2011) 2007~2008, 276 days of data



## Select muon track events from -5° to 85° in declination.



 $\begin{array}{l} \text{Measure } \Delta N = (N^{\text{bkg}}{}_{\text{on}} + N^{\text{sig}}{}_{\text{on}}) - (N^{\text{bkg}}{}_{\text{off}} + N^{\text{sig}}{}_{\text{off}}) \\ \approx \Delta N^{\text{sig}} \end{array}$ 





#### This result is almost independent of galactic halo model



PAMELA: GeV positron excess Fermi: electron spectra

#### Constraints for annihilation to $\mu^+\mu^-$



We are interested in low DM mass case Consider 10 GeV energy threshold for track and shower

#### Atmospheric neutrino fluxes averaged for $-1 \leq \cos \xi \leq -0.4$

F. F. Lee, G. L. Lin, Astropart. Phys. 25, 2006









#### **Comparison with Honda flux and AMANDA-II measurement**

## Neutrino flux from DM decay in the galactic halo

$$\frac{\mathrm{d}\Phi_{\nu_i}}{\mathrm{d}E_{\nu_i}} = \frac{\Delta\Omega}{4\pi} \frac{1}{m_{\chi}\tau_{\chi}} \left( \sum_F B_F \frac{\mathrm{d}N_{\nu_i}^F}{\mathrm{d}E} \right) R_{\oplus} \rho_{\oplus} \times J_1(\Delta\Omega) \quad \propto \frac{\rho}{m_{\chi}\tau_{\chi}}$$

- $R_{\oplus} = 8.5 \text{kpc}$  distance from the galactic center to the solar system  $\rho_{\oplus} = 0.3 \text{ GeV/cm}^3$ : DM density in the solar neighborhood  $dN_{v_i}^F/dE$ : neutrino spectrum per decay for a given decay channel *F*  $\tau_{\chi}$ : DM lifetime
- $J_1(\Delta \Omega)$  is the DM distribution integrated over the line-of-sight (l.o.s) for decay and averaged over a solid angle  $\Delta \Omega = 2\pi(1 \cos\psi_{max})$

$$J_{1}(\Delta \Omega) = \frac{1}{\Delta \Omega} \int_{\Delta \Omega} d\Omega \int_{l.o.s} \frac{dl}{R_{\oplus}} \left( \frac{\rho(r(l,\psi))}{\rho_{\oplus}} \right)^{1}$$

Navarro-Frenk-White (NFW) DM density profile

$$\rho(r) = \rho_s \left(\frac{R_s}{r}\right) \left(\frac{R_s}{R_s + r}\right)^2$$



## Neutrino flux from DM annihilation in the galactic halo

$$\frac{\mathrm{d}\Phi_{v_i}}{\mathrm{d}E_{v_i}} = \frac{\Delta\Omega}{4\pi} \frac{\langle \sigma \upsilon \rangle}{2m_{\chi}^2} \left( \sum_F B_F \frac{\mathrm{d}N_{v_i}^F}{\mathrm{d}E} \right) R_{\oplus} \rho_{\oplus}^2 \times J_2(\Delta\Omega) \propto \frac{\rho^2 \langle \sigma \upsilon \rangle}{2m_{\chi}^2}$$

•  $<\sigma v>$  is the thermally averaged annihilation cross section

 $\langle \sigma \upsilon \rangle = B \langle \sigma \upsilon \rangle_0$ 

**B** : boost factor  $. < \sigma v >_0 = 3 \times 10^{-26} \text{ cm}^3 \text{s}^{-1}$ : typical cross section for DM relic density.

•  $J_2(\Delta\Omega)$  is the line-of-sight (l.o.s) integral for annihilation

$$J_{2}(\Delta \Omega) = \frac{1}{\Delta \Omega} \int_{\Delta \Omega} d\Omega \int_{l.o.s} \frac{dl}{R_{\oplus}} \left( \frac{\rho(r(l,\psi))}{\rho_{\oplus}} \right)^{2}$$

## **Neutrino fluxes on Earth**

$$\begin{pmatrix} \Phi_{\nu_e} \\ \Phi_{\nu_{\mu}} \\ \Phi_{\nu_{\tau}} \end{pmatrix} = \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ P_{\mu e} & P_{\mu\mu} & P_{\mu\tau} \\ P_{\tau e} & P_{\tau\mu} & P_{\tau\tau} \end{pmatrix} \begin{pmatrix} \Phi_{\nu_e}^0 \\ \Phi_{\nu_{\mu}}^0 \\ \Phi_{\nu_{\tau}}^0 \end{pmatrix} = P \begin{pmatrix} \Phi_{\nu_e}^0 \\ \Phi_{\nu_{\mu}}^0 \\ \Phi_{\nu_{\tau}}^0 \end{pmatrix}$$

J. G. Learned, Astropart. Phys. 3, 1995 H. Athar et al. Phys. Rev. D62, 2000 L. Bento et al. Phys. Lett. B476, 2000 K. C. Lai et al. Phys. Rev. D82, 2010

 $\left\{ \begin{array}{l} \Phi^0_{\nu_{\alpha}} : \text{neutrino flux at the astrophysical source} \\ \Phi_{\nu_{\alpha}} : \text{neutrino flux measured on the Earth} \\ P_{\alpha\beta} : \text{probability of the oscillation} \quad \nu_{\beta} \rightarrow \nu_{\alpha} \end{array} \right.$ 

In the tribimaximal limit of neutrino mixing angles:

 $\sin^2 \theta_{23} = 1/2, \sin^2 \theta_{12} = 1/3, \sin^2 \theta_{13} = 0$ 

$$P = \begin{pmatrix} \frac{5}{9} & \frac{2}{9} & \frac{2}{9} \\ \frac{2}{9} & \frac{7}{18} & \frac{7}{18} \\ \frac{2}{9} & \frac{7}{18} & \frac{7}{18} \\ \frac{2}{9} & \frac{7}{18} & \frac{7}{18} \end{pmatrix} \qquad \Longrightarrow \qquad \begin{cases} \Phi_{\nu_e} = \frac{5}{9} \Phi_{\nu_e}^0 + \frac{2}{9} \Phi_{\nu_{\mu}}^0 + \frac{2}{9} \Phi_{\nu_{\tau}}^0 \\ \Phi_{\nu_{\mu}} = \Phi_{\nu_{\tau}} = \frac{2}{9} \Phi_{\nu_{e}}^0 + \frac{7}{18} \Phi_{\nu_{\mu}}^0 + \frac{7}{18} \Phi_{\nu_{\tau}}^0 \end{cases}$$

• with recently measured  $\theta_{13}$  value:

 $\sin^2 2\theta_{13} = 0.092$ 

Dayabay best fit, PRL108, 171803(2012) See also Double Chooz, PRL108, 131801 (2011) RENO arXiv:1204.0626

 $P = \begin{pmatrix} 0.53 & 0.26 & 0.21 \\ 0.26 & 0.37 & 0.37 \\ 0.21 & 0.37 & 0.42 \end{pmatrix}$ 

#### **Event Rates**

$$\Gamma_{\text{track}} = \int_{E_{\mu}^{\text{th}}}^{E_{\text{max}}} dE_{\mu} \int_{E_{\mu}}^{E_{\text{max}}} dE_{\nu_{\mu}} N_A \rho_{\text{ice}} V_{\text{tr}} \times \frac{d\Phi_{\nu_{\mu}}}{dE_{\nu_{\mu}}} \cdot \frac{d\sigma_{\nu_{\mu}N}^{CC} (E_{\nu_{\mu}}, E_{\mu})}{dE_{\mu}} + (\nu \to \overline{\nu})$$

$$\Gamma_{\underline{\text{cascade}}} = \int_{E_{\text{shower}}}^{E_{\text{max}}} dE_{\text{shower}} \int_{E_{\text{Shower}}}^{E_{\text{max}}} dE_{\nu} N_{A} \rho_{\text{ice}} V_{\text{casc}} \times \frac{d\Phi_{\nu}}{dE_{\nu}} \cdot \frac{d\sigma_{\nu N}(E_{\nu}, E_{\text{shower}})}{dE_{\text{shower}}} + (\nu \rightarrow \overline{\nu})$$

$$\rightarrow (\nu_{e}N)_{\text{CC}} , (\nu_{e}N)_{\text{NC}} , (\nu_{\mu}N)_{\text{NC}} , (\nu_{\tau}N)_{\text{CC}} , (\nu_{\tau}N)_{\text{NC}}$$

- $\rho_{ice} = 0.9 \text{ g cm}^{-3}$  is the density of ice;  $N_A = 6.022 \times 10^{23} \text{ g}^{-1}$  is Avogadro's number
- $V_{tr} \approx 0.04 \text{ km}^3$  is the effective volume of IceCube DeepCore array for muon track events  $V_{casc} \approx 0.02 \text{ km}^3$  is the effective volume of IceCube DeepCore array for cascade events
- $E_{\text{max}}$  is taken as  $m_{\chi}$  for DM annihilatin;  $E_{\text{max}}$  is taken as  $\frac{m_{\chi}}{2}$  for DM decay
- $\frac{\mathrm{d}\Phi_{\nu_e}}{\mathrm{d}E_{\nu_e}}$  is taken from M. Honda et al. Phys. Rev. D75, 2007

#### **Constraints for DM annihilation cross section**



50 GeV threshold from A. E. Erkoca, M. H. Reno, I. Sarcevic, Phys. Rev. D82, 2010



# **Constraints on various modes**

## **Constraints for DM decay** $\chi \rightarrow \mu^+ \mu$ **track & cascade**)



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## Compare with Fermi-LAT constraints Detecting extragalactic isotropic gamma-ray







20 MeV to 300 GeV

## **Compare with Fermi-LAT constraints**



More comparisons



## More comparisons



# **Summary**

(1) We employ NFW DM profile to calculate the track and cascade event rates in IceCube DeepCore due to neutrino fluxes from WIMP annihilations and decays in the galactic halo.

(2) We take into account neutrino oscillations and calculate the event rates due to atmospheric neutrino background.

(3) Cascade events provide stronger constraints on DM annihilation cross section and DM decay time than the corresponding constraints provided by track events with the same threshold energy.

(4) Fermi\_LAT gamma ray constraints are generally more stringent than those expected at IceCube DeepCore. On the other hand, DM annihilating directly into neutrino pair is not constrained by gamma ray.

# Summary

(5) KM3NeT situated in northern hemisphere could provide better sensitivity on DM annihilation cross section.

http://www.km3net.org/TDR/TDRKM3NeT.pdf