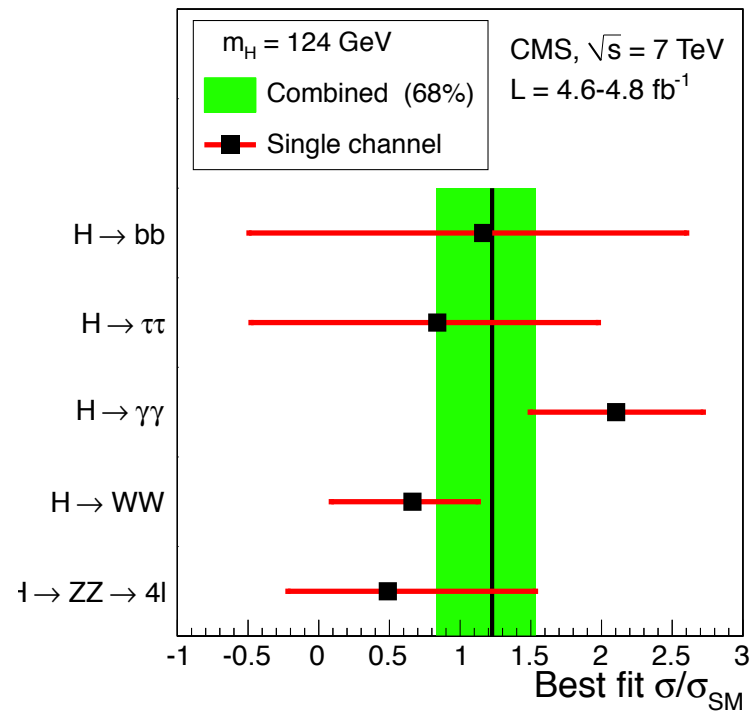
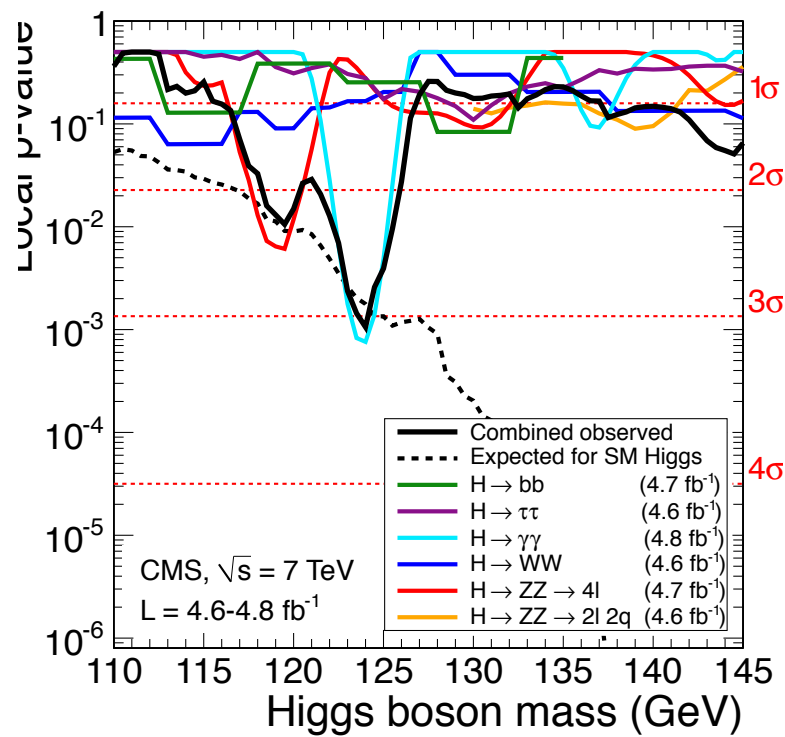


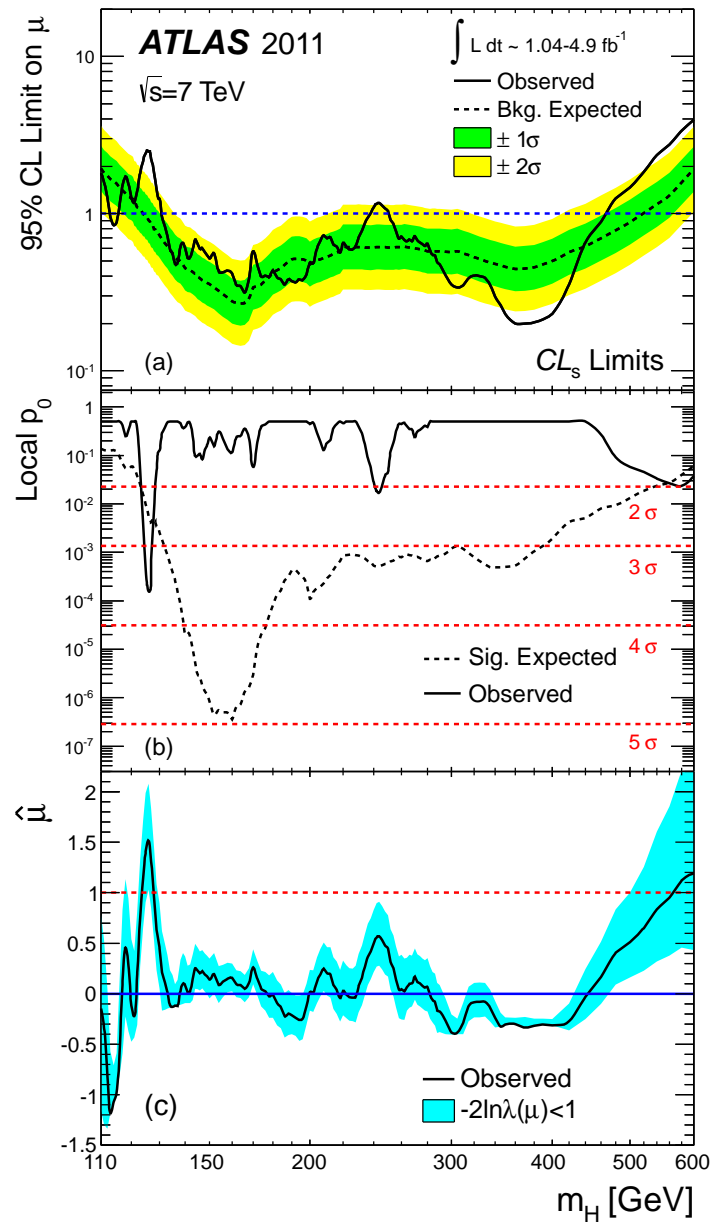
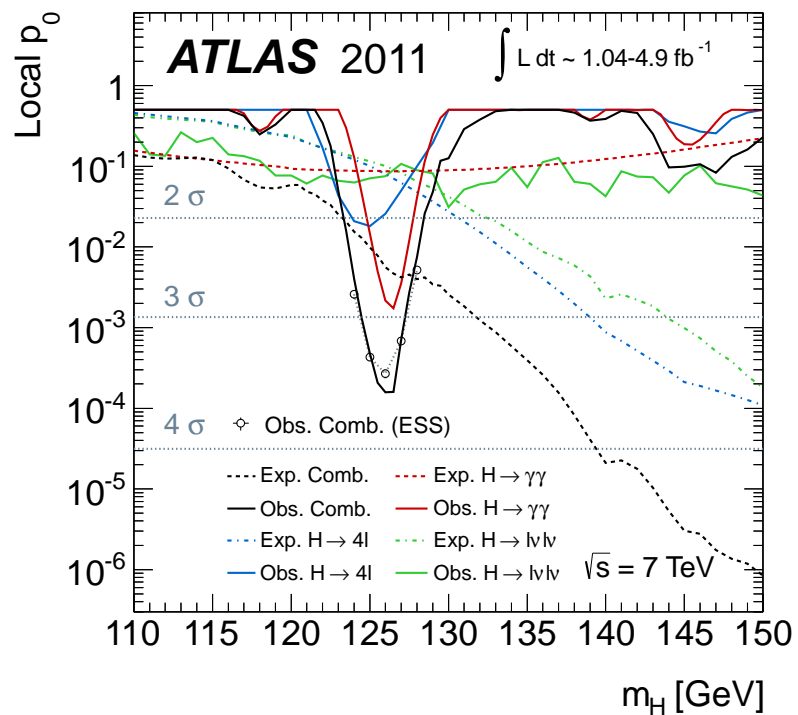
# Higgs Bosons and Z' Phenomenology of $U(1)$ -MSSM

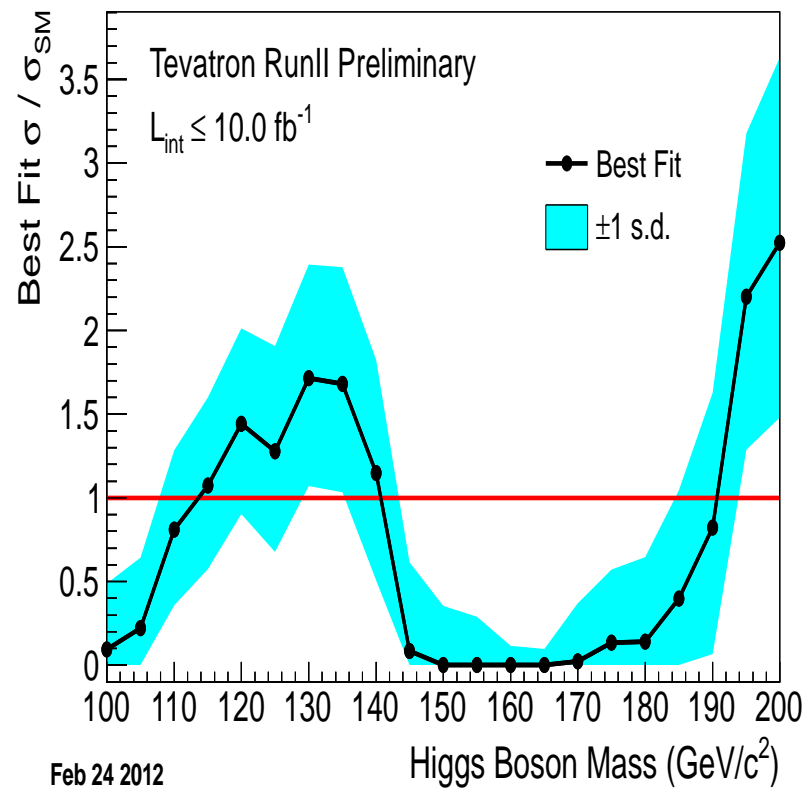
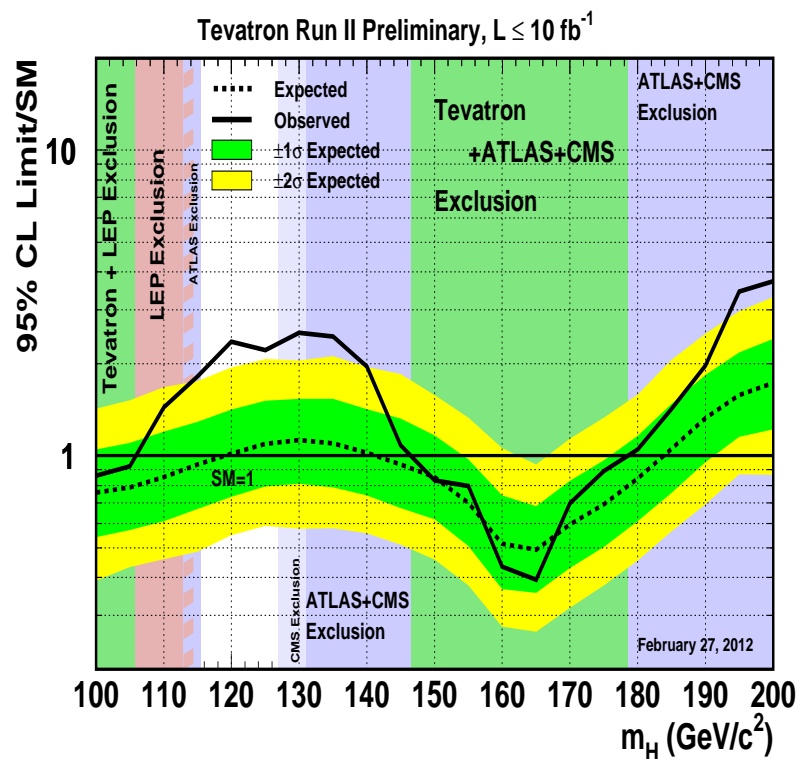
KC, Chun-Fu Chang, Yi-Chuan Lin, TC Yuan (arXiv:1202.0054 [hep-ph])

Konkuk WCU/NTHU

Cross Strait Meeting 2012







## Various Interpretations

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- MSSM: the most straight forward, SUSY predicts a light Higgs boson. But to give a 125 GeV Higgs puts a tight constraint on the stop mass sector, and difficult to enhance the  $\gamma\gamma$  rate. (Baer et al; Heinemeyer et al; Arbey et al; Draper et al; Carena et al; Cao et al, ...)
- NMSSM: easier to obtain a 125 GeV Higgs boson, but still controversial to obtain enhanced  $\gamma\gamma$  rate. (Gunion et al; Ellwanger; King et al., )
- RS Radion/dilaton: the anomaly couplings to  $gg$  and  $\gamma\gamma$  easily enhance the diphoton rate. (KC and TC, PRL)
- 2HDM: Ferreira et al.
- 4th generation: Guo, Ren, and He
- Inert Higgs doublet model: Arhrib et al.
- $U(1)$ -extended MSSM: C.F. Chang, KC, Y.C. Lin, TC

## U(1)-MSSM

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- Extra  $U(1)$ s exist in the breaking of  $E_6$ :

$$E_6 \rightarrow SO(10) \times U(1)_\psi, \quad SO(10) \rightarrow SU(5) \times U(1)_\chi.$$

and each family of left-handed fermions is promoted to a fundamental **27**. The two  $U(1)$ s can mix to form various  $Z'$  models.

- We require one of the  $U(1)$  remains unbroken until TeV and then broken to give mass to  $Z'$  and  $Z'$ -ino.

$SO(16)$	$SU(5)$	$2\sqrt{10}Q'_\chi$	$2\sqrt{6}Q'_\psi$	$2\sqrt{15}Q'_\eta$
<b>16</b>	<b>10</b> ( $u, d, u^c, e^c$ )	-1	1	-2
	<b>5*</b> ( $d^c, \nu, e^-$ )	3	1	1
	$\nu^c$	-5	1	-5
<b>10</b>	<b>5</b> ( $D, H_u$ )	2	-2	4
	<b>5*</b> ( $D^c, H_d$ )	-2	-2	1
<b>1</b>	<b>1S</b>	0	4	-5

- Based on the  $U(1)$  quantum number the allowed terms in superpotential

$$W_{\text{eff}} = \epsilon_{ab} \left[ y_{ij}^u Q_j^a H_u^b U_i^c - y_{ij}^d Q_j^a H_d^b D_i^c - y_{ij}^l L_j^a H_d^b E_i^c + h_s S H_u^a H_d^b \right]$$

The  $U(1)$  charges of  $S, H_u, H_d$  are related by  $Q'_{H_u} + Q'_{H_d} + Q'_S = 0$ .

- The gauge interactions involving  $Z'$  and the supermultiplet  $(\phi, \psi)$  is

$$\mathcal{L} = \frac{1}{2} \bar{\psi}_i i\gamma^\mu D_\mu \psi_i + (D^\mu \phi_i)^\dagger (D_\mu \phi_i),$$

where

$$D_\mu \phi_i = \left[ \partial_\mu + ieQA_\mu + i\frac{g}{\sqrt{2}}(\tau^+ W_\mu^+ + \tau^- W_\mu^-) + ig_1(T_{3L} - Qx_w)Z_\mu + ig_2 Z'_\mu Q' \right] \phi_i.$$

- The  $Z'$  mass arises from VEVs of  $H_u, H_d, S$ :

$$m_{Z'} \approx g_2 (Q_{H_u}'^2 v_u^2 + Q_{H_d}'^2 v_d^2 + Q_s'^2 v_s^2)^{1/2}$$

The  $Z'$ -ino mass comes from the soft SUSY breaking gaugino mass term.

## Physical States

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- Higgs bosons:

$$H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix} \quad \text{and} \quad S.$$

Similar to NMSSM, UMSSM has 3 CP-even Higgs bosons, 1 charged Higgs, but **only one pseudoscalar boson, because one is eaten by the  $Z'$  boson.**

$$\begin{aligned} H_d^0 &= \frac{1}{\sqrt{2}} (v_d + \phi_d + i\chi_d), \\ H_u^0 &= \frac{1}{\sqrt{2}} (v_u + \phi_u + i\chi_u), \\ H_s^0 &= \frac{1}{\sqrt{2}} (v_s + \phi_s + i\chi_s), \end{aligned}$$

So  $h_1, h_2, h_3, A, H^\pm$ .



- Neutralinos, total 6: in the basis of  $(\tilde{B}, \tilde{W}^3, \tilde{h}_d^0, \tilde{h}_u^0, \tilde{S}, \tilde{Z}')^T$

$$\mathcal{M}_N = \left( \begin{array}{cccc|cc} M_1 & 0 & -\frac{e}{2c_w} v_d & \frac{e}{2c_w} v_u & 0 & 0 \\ 0 & M_2 & \frac{g}{2} v_d & -\frac{g}{2} v_u & 0 & 0 \\ -\frac{e}{2c_w} v_d & \frac{g}{2} v_d & 0 & -\mu_{\text{eff}} & -\frac{h_s}{\sqrt{2}} v_u & g_2 Q'_{H_d} v_d \\ \frac{e}{2c_w} v_u & -\frac{g}{2} v_u & -\mu_{\text{eff}} & 0 & -\frac{h_s}{\sqrt{2}} v_d & g_2 Q'_{H_u} v_u \\ \hline 0 & 0 & -\frac{h_s}{\sqrt{2}} v_u & -\frac{h_s}{\sqrt{2}} v_d & 0 & g_2 Q'_S v_s \\ 0 & 0 & g_2 Q'_{H_d} v_d & g_2 Q'_{H_u} v_u & g_2 Q'_S v_s & M_{\tilde{Z}'} \end{array} \right) .$$

## Two Higgs Scenarios

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- **The SM-like Higgs boson:  $130 < M_{h_{\text{SM-like}}} < 141$  GeV.** We have to hide it from showing up in the  $WW$ ,  $ZZ$ ,  $\gamma\gamma$ ,  $b\bar{b}$ , and  $\tau^+\tau^-$  channels. A possible way is to allow decay into neutralinos.
- **The SM-like Higgs boson:  $120 < M_{h_{\text{SM-like}}} < 130$  GeV.** This is what LHC and Tevatron are seeing. So have to make sure that the  $WW$ ,  $ZZ$ ,  $\gamma\gamma$ ,  $b\bar{b}$ , and  $\tau^+\tau^-$  channels are there. We have to suppress decay into neutralinos.

## Scanning of Parameter Space

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- Fix some parameters:

$$M_1 = 0.5M_2 = 0.2 \text{ TeV}, \quad M_3 = 2 \text{ TeV};$$

$$M_{\tilde{Q}} = M_{\tilde{U}} = A_t = 1 \text{ TeV}, \quad M_{\tilde{L}} = M_{\tilde{E}} = 0.2 \text{ TeV} . M_S = 0.5 \text{ TeV}, \quad A_s = 0.5 \text{ TeV}$$

- Vary

$$0.2 \text{ TeV} < v_s < 2 \text{ TeV}, \quad 0.2 < h_s < 0.7, \quad 1.1 < \tan \beta < 40, \quad 0.2 \text{ TeV} < M_{\tilde{Z}'} < 2 \text{ TeV}$$

- Constraints:

1. **Chargino mass > 94 GeV**
2. Invisible width of  $Z$

$$\Gamma(Z \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0) = \frac{g_1^2}{96\pi} \left( |N_{13}|^2 - |N_{14}|^2 \right)^2 m_Z \left( 1 - \frac{4m_{\tilde{\chi}_1^0}^2}{m_Z^2} \right)^{3/2} < \mathbf{3 \text{ MeV}}$$

## Higgs Decays

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- Higgs decays into gauge bosons
- Higgs decays into fermions
- Higgs decays into neutralinos
- Higgs decays into lighter Higgs

Scenario 1:  $130 < M_{h_{\text{SM-like}}} < 141$  GeV.

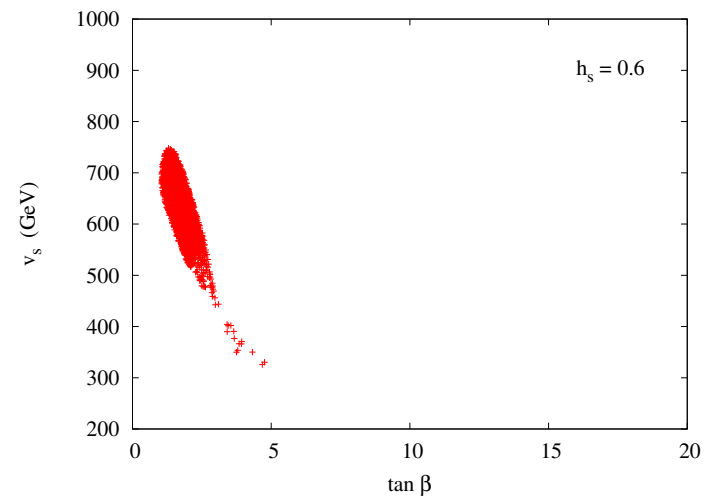
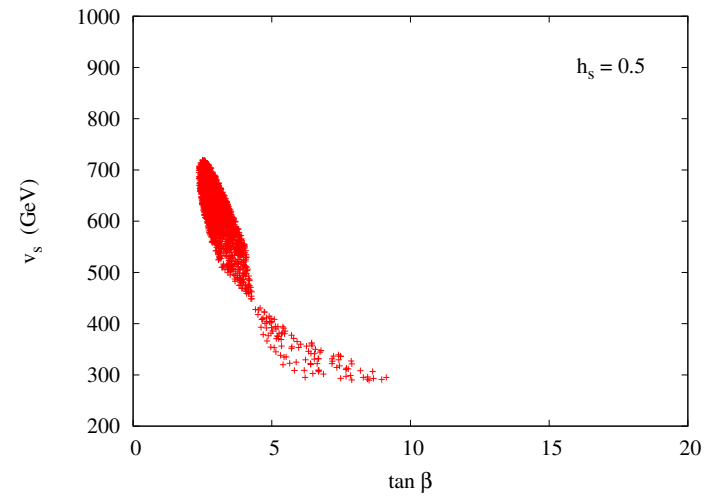
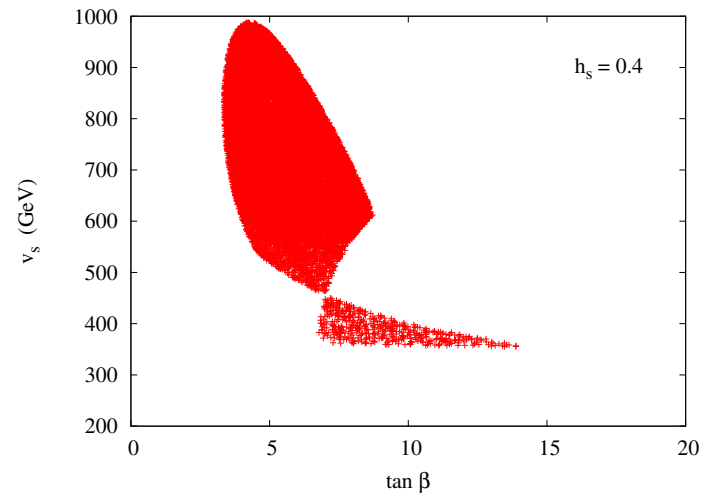
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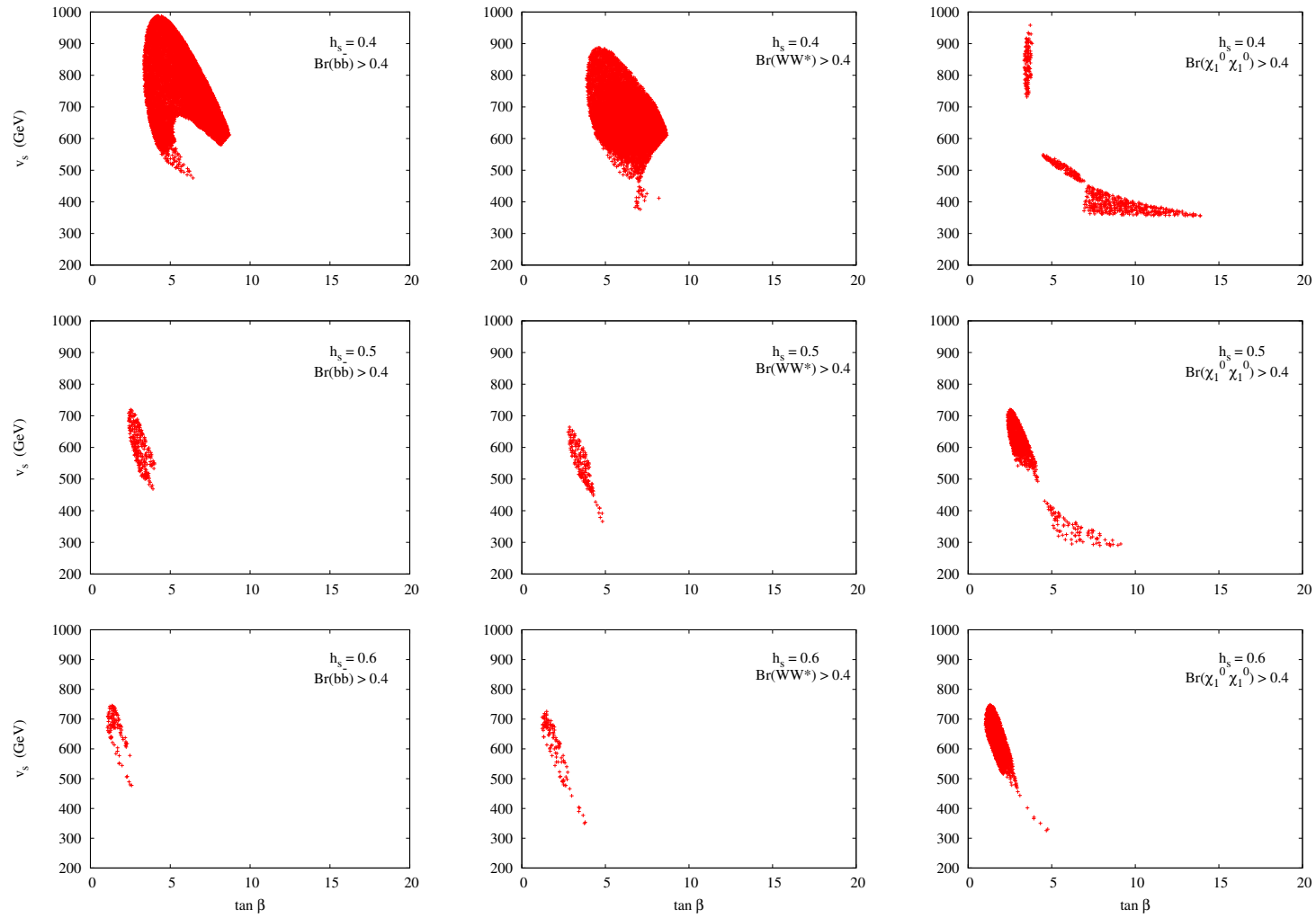
- The physical Higgs states are

$$h_k = O_{k1}\phi_d + O_{k2}\phi_u + O_{k3}\phi_s$$

Require the singlet component to be small:  $O_{k3}^2 < 0.1$  to be the SM-like.

- In our scan we do not find more than one SM-like Higgs boson.
- A smaller  $h_s$  is more likely to give a Higgs boson in mass range 130 – 140 GeV.
- If  $B(h_{\text{SM-like}} \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0) > 0.4$ , we are able to hide the Higgs boson under the current data. Larger  $h_s$  will make invisible mode dominant.



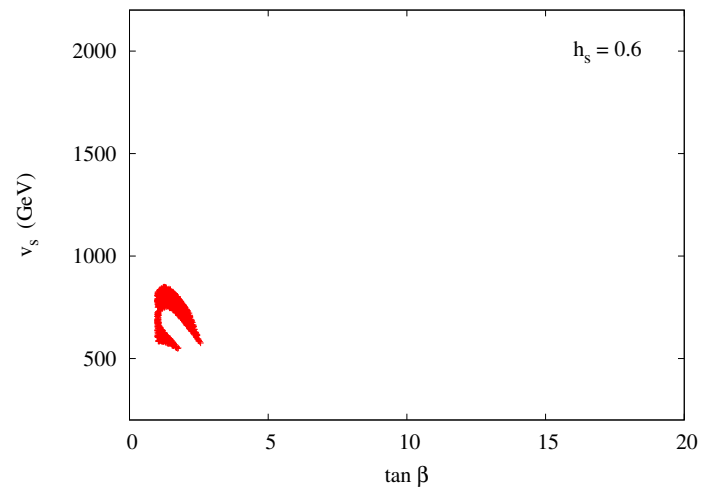
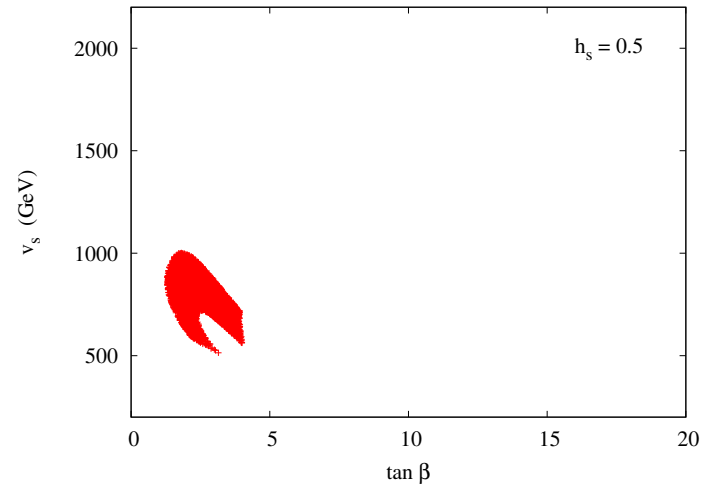
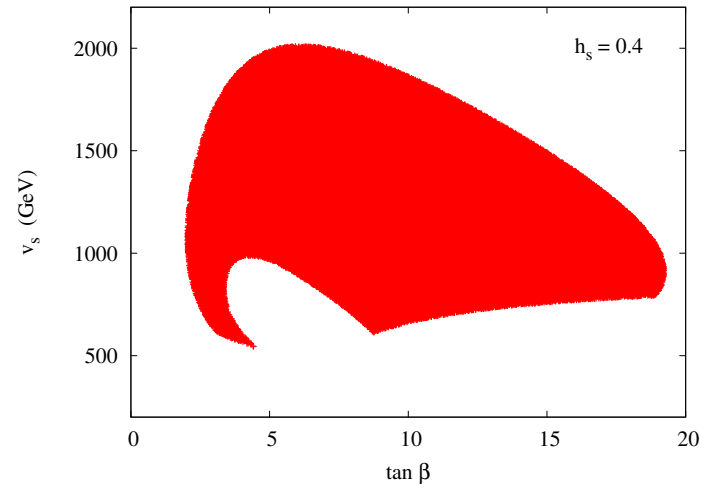


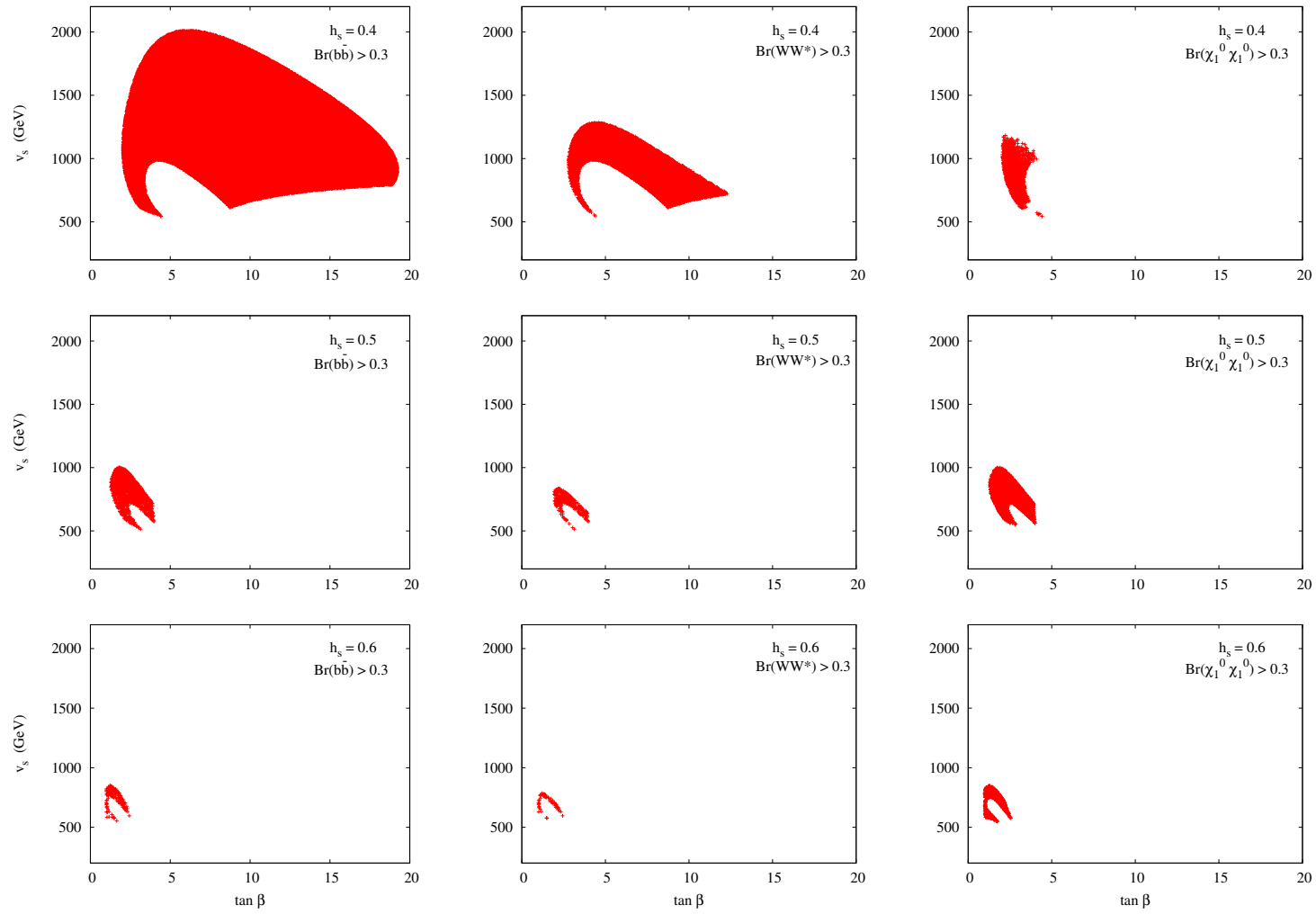
Scenario 2:  $120 < M_{h_{\text{SM-like}}} < 130 \text{ GeV}$ .

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- A smaller  $h_s$  is more likely to give  $120 \text{ GeV} < M_{\text{SM-like}} < 130 \text{ GeV}$ .
- Again larger  $h_s$  will make invisible mode dominant.







## Conclusions

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- Extensions to the MSSM easily raises the Higgs boson mass, e.g., NMSSM, UMSSM.
- The UMSSM only has one pseudoscalar such that  $h_{\text{SM-like}}$  cannot decay into two  $a_1$ , in contrast to NMSSM.
- If the  $h_{\text{SM-like}}$  falls in the range 130 – 141 GeV, it can be hidden from current data by decaying into neutralinos. Still allowed.
- If the  $h_{\text{SM-like}}$  falls in the range 120 – 130 GeV, it can decay similar to the SM Higgs boson. Thus, allowed by current data.