### Heavy Charged Higgs@LHC

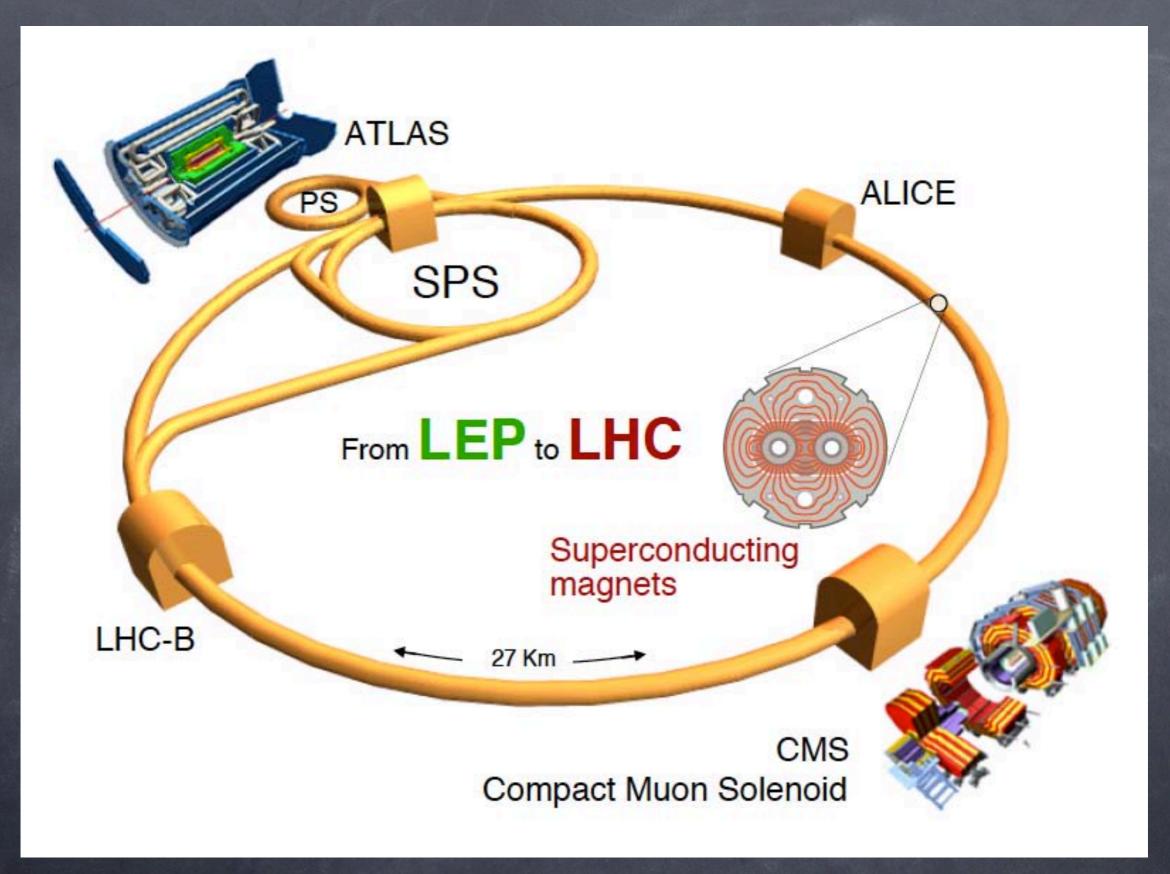
Qi-Shu Yan GUCAS

In collaboration with Shuo Yang (Peking U & Dalian U) arXiv: 1111.4530;Published in JHEP1202(2012)074

Cross Strait Meeting, ChongQing, May 7-12, 2012

### Outline

1. Why A Heavy Charged Higgs Boson?
2. What is jet/Substructure/top taggers?
3. Feasibility at the LHC with a Hybrid-R method
4. Conclusions

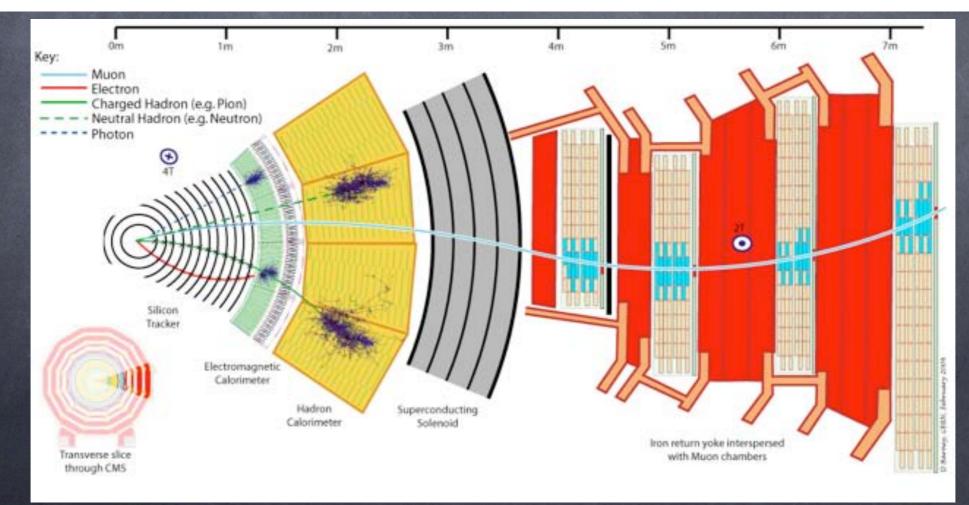


#### Collecting data 7TeV~5/fb, 8TeV~1/fb

#### j (Jet) Cluster in EM and hadronic calorimeters (and inner tracker) $\gamma$ Photon EM cluster without matching track

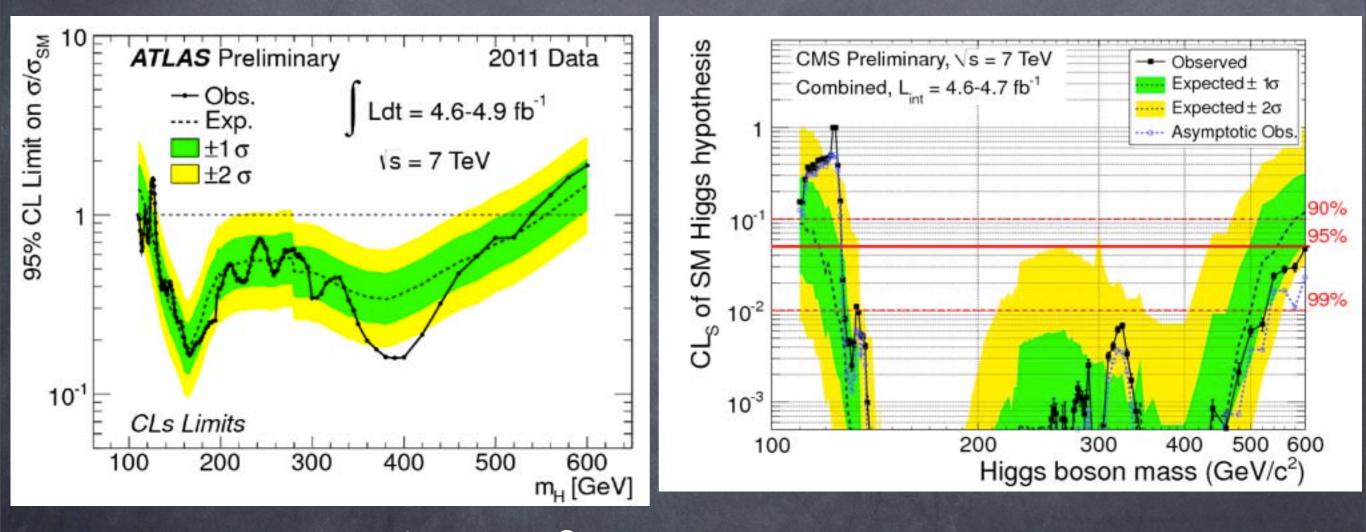
- e Electron EM cluster with matching track
- μ Muon Matching tracks in inner and muon trackers
- τ Tau lepton Narrow jet with matching track(s)

MET Missing  $E_T$  pT required to balance all of the above (and more)



#### CMS Detector as an example

# 1. Why A Heavy Charged Higgs Boson?

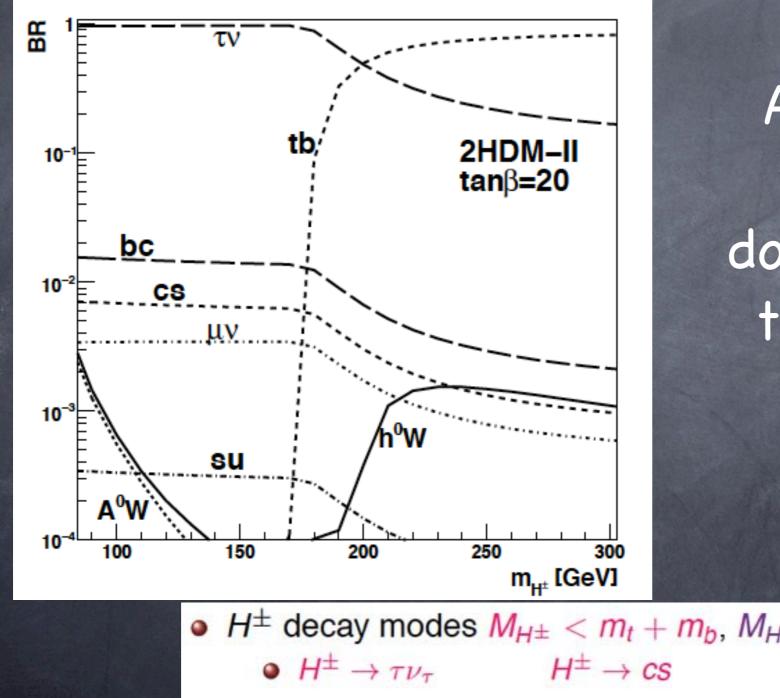


There are evidence for a light Higgs boson around 125 GeV. Whether the Higgs is SM-like? What's is the Higgs Sector?

# 1. Why Charged Higgs Boson(s)?

- Charge Higgs boson(s) are predictions of most of extension of the SM, i.e. 2HDM, MSSM, mutli-Higgs model, little Higgs models.
- Discovery charged Higgs boson(s) is the herald of New Physics.
- Discovery/Ruling-out of charged Higgs boson(s) is helpful to pinpoint the Higgs sector (EW symmetry breaking).

## 1. Why Heavy Charged Higgs Boson?



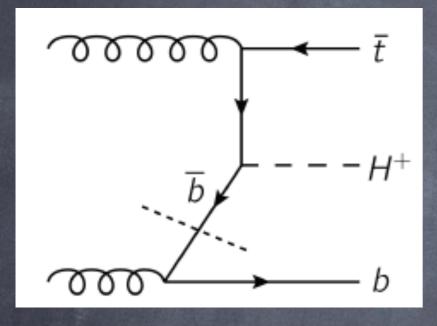
A Heavy Charged Higgs can dominantly decay to t+b (i.e. 2HDM-II) if m>800 GeV

•  $H^{\pm}$  decay modes  $M_{H^{\pm}} < m_t + m_b, M_{H^{\pm}} > m_t + m_b$ : •  $H^{\pm} \to \tau \nu_{\tau}$   $H^{\pm} \to cs$   $H^{\pm} \to t^*b \to W^{\pm}\bar{b}b$  $H^{\pm} \to A^0 W^{\pm}$   $H^{\pm} \to h^0 W^{\pm}$   $H^{\pm} \to tb$ 

# 1. Why A Heavy Charged Higgs Boson? Charged Higgs boson search

- For  $m_{H^+} < m_t$ , search for  $t \to H^+ b$ in top pair decay. Tevatron set a limit for charged Higgs mass up to ~ 160.
- At LHC, the main production mechanism gb → tH<sup>-</sup>. Previous studies for charged Higgs search mainly forcus on 300-600GeV gb → tH<sup>-</sup> → tτν̄ gb → tH<sup>-</sup> → tt̄b → bqqτ(hadronic)νbb gb → tH<sup>-</sup> → tt̄b → bqqbℓνb

# 1. Why A Heavy Charged Higgs Boson?What's New?



Difficulties in previous study:

- Combinatorics is huge
- Lack of powerful observables to suppress background

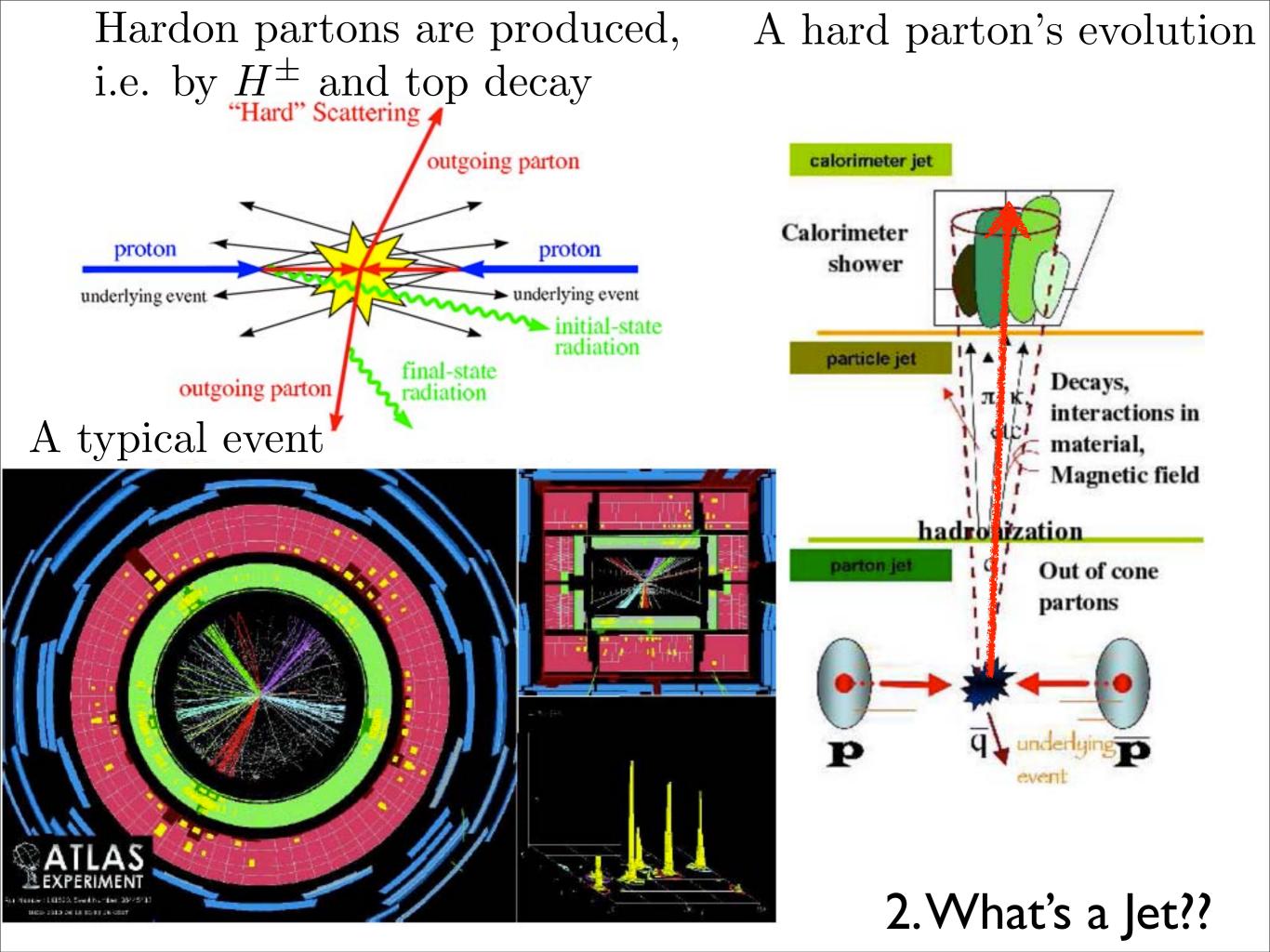
 $pp \rightarrow tH^- \rightarrow t\bar{t}b$   $pp \rightarrow tH^-b \rightarrow t\bar{t}bb$ Test Field: the full hadronic channel

New Things 800-1500GeV
Combinatorics can be greatly reduced if the charged Higgs boson is heavy
Powerful observables can be found after the full reconstruction Although there are 7 or more objects in the final states (at parton level for our signals)

- The top from heavy charged Higgs boson decay will be captured by using top tagger(s). A larger cone
- The b quark from heavy charged Higgs boson is more energetic than jets from the associated top quark.
- The rest of jets will be used to reconstruct the associated top quark

#### A smaller cone

Combinatoric issue can be overcome and mass bumps can be reconstructed.

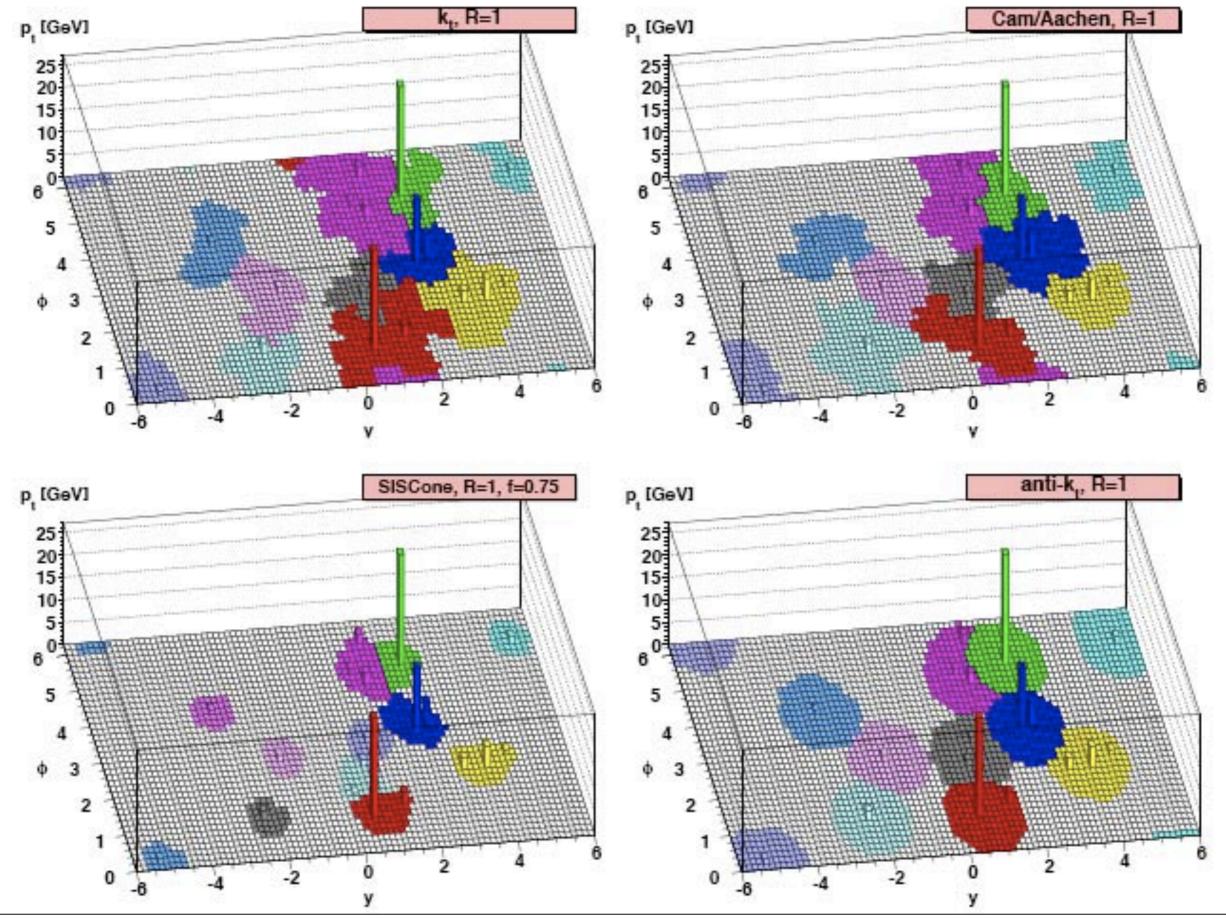


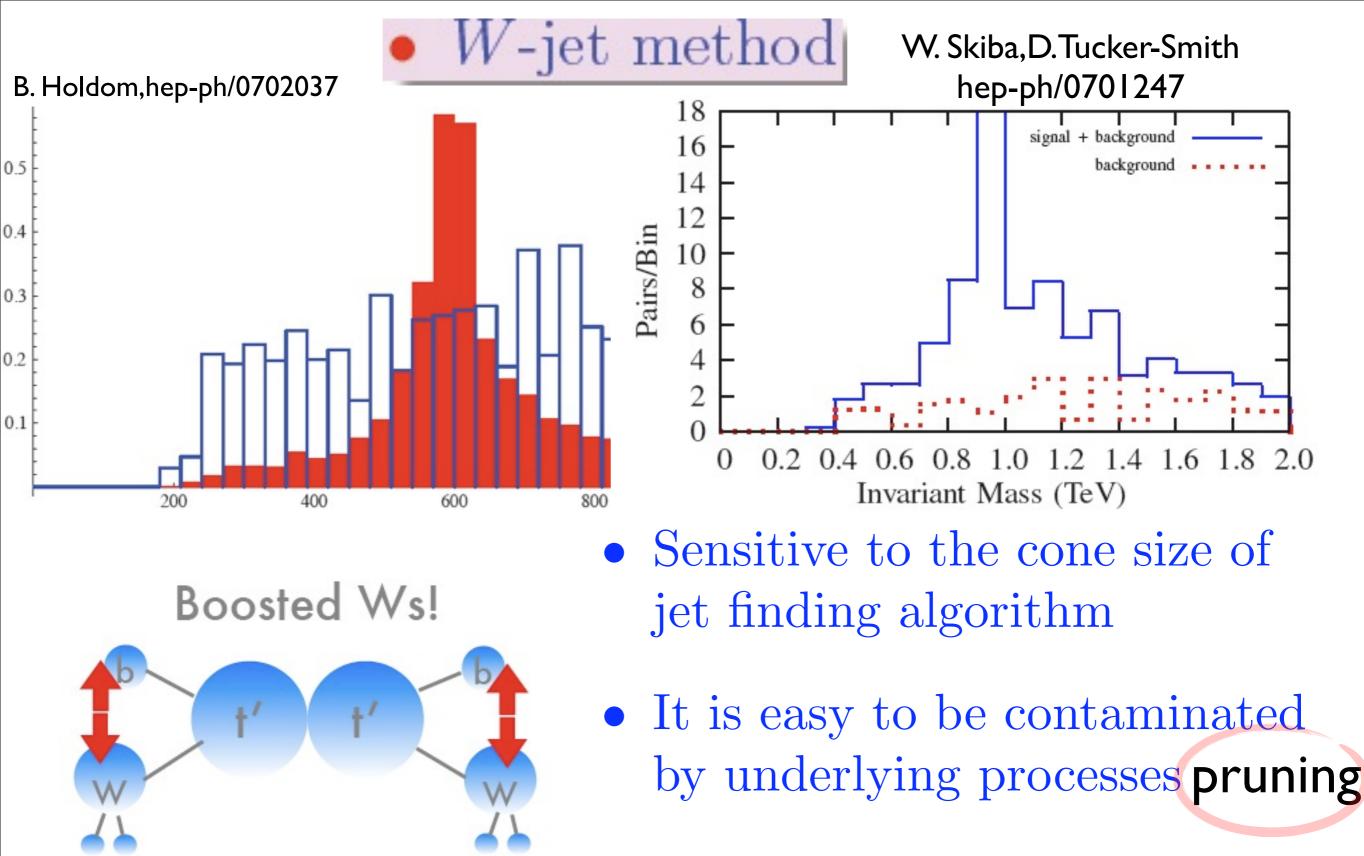
Compute the smallest "distance"  $d_{ij}$  or  $d_{iB}$  and either cluster i and j together or identify i as a jet

$$d_{ij} = \min\{k_{Ti}^{p}, k_{Tj}^{p}\}\Delta R_{ij}/R, \qquad d_{iB} = k_{Ti}^{p}$$
$$\Delta R^{2} = \Delta \eta^{2} + \Delta \phi^{2}$$

Algorithm	p	clusters first	comment
k <sub>T</sub> /Durham	> 0	softest	leads to very irregular jets includes a lot of underlying event hard to get jet energy scale right
Cambridge/Aachen	= 0	closest	still leads to very irregular jets similar problems to kT algorithm
anti-k <sub>T</sub>	< 0	hardest	shape of jet insensitive to soft particles √ cone-like jets √ may be easier to get jet energy scale right √

#### Towards Jetography: 0906.1833, G.P. Salam



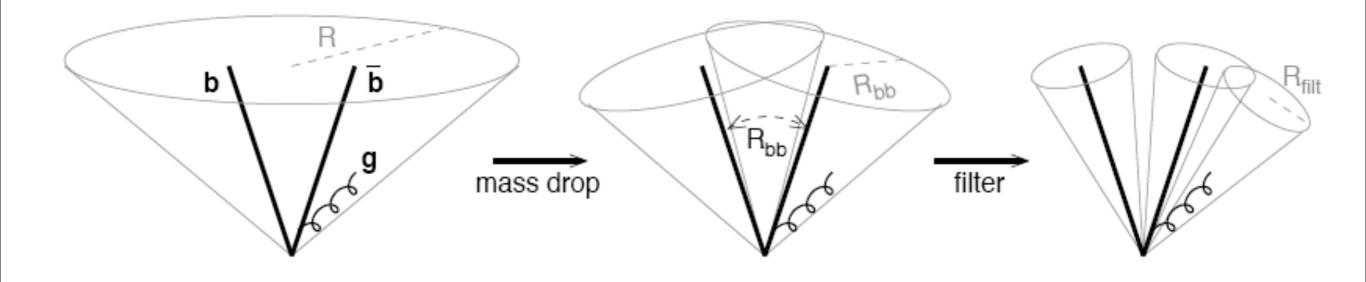


2. Jet mass/sub.

• W must be highly boosted need larger mass of heavy quarks  $m_{t'} > 500 \text{GeV}$ 

Jet Substructure study entered into a new era.

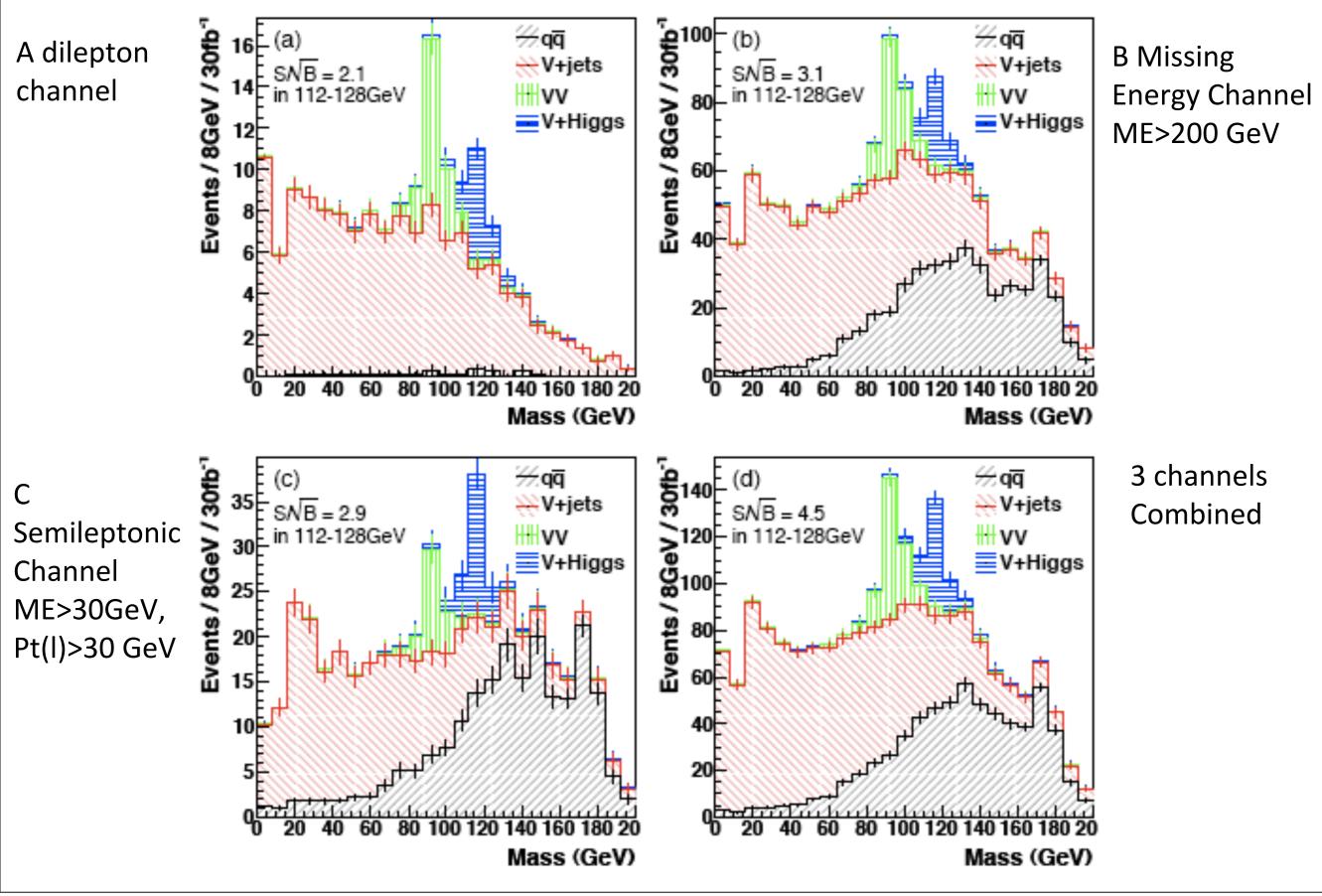
Higgs tagger (BDRS): 0802.2470 J.M. Butterworth, A.R. Davison, M.Rubin, G.P. Salam



#### Optimized for R =1.2

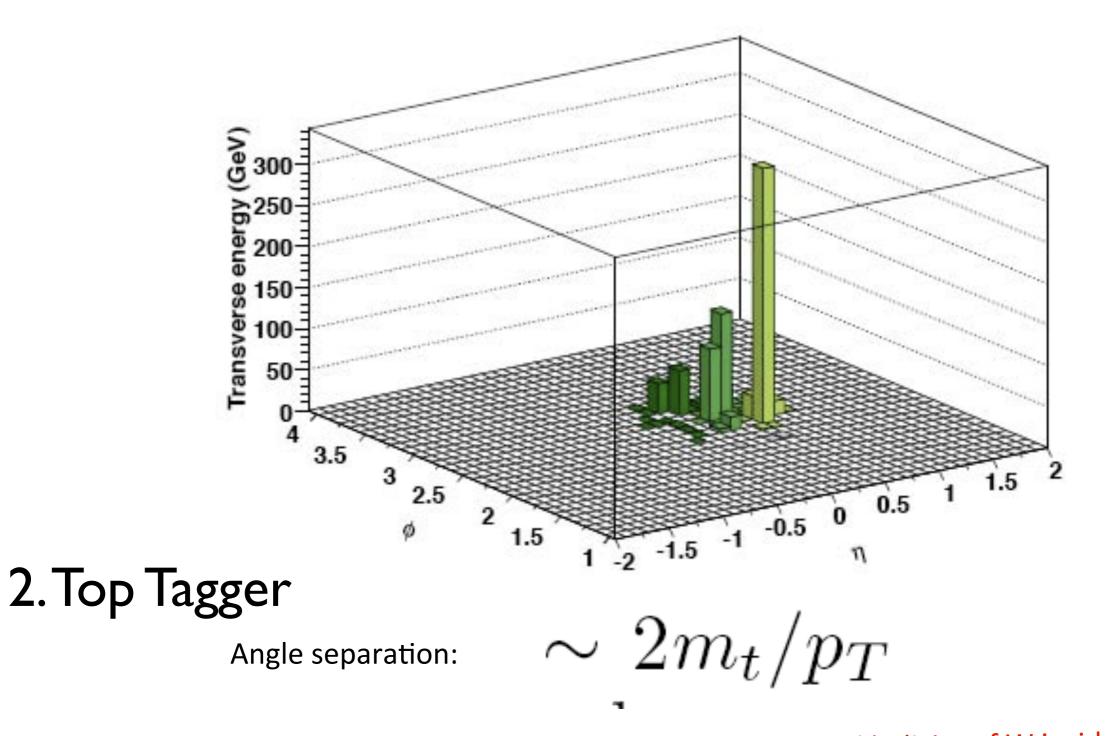
#### 2. Higgs tagger

Higgs tagger (BDRS): 0802.2470 J.M. Butterworth, A.R. Davison, M.Rubin, G.P. Salam



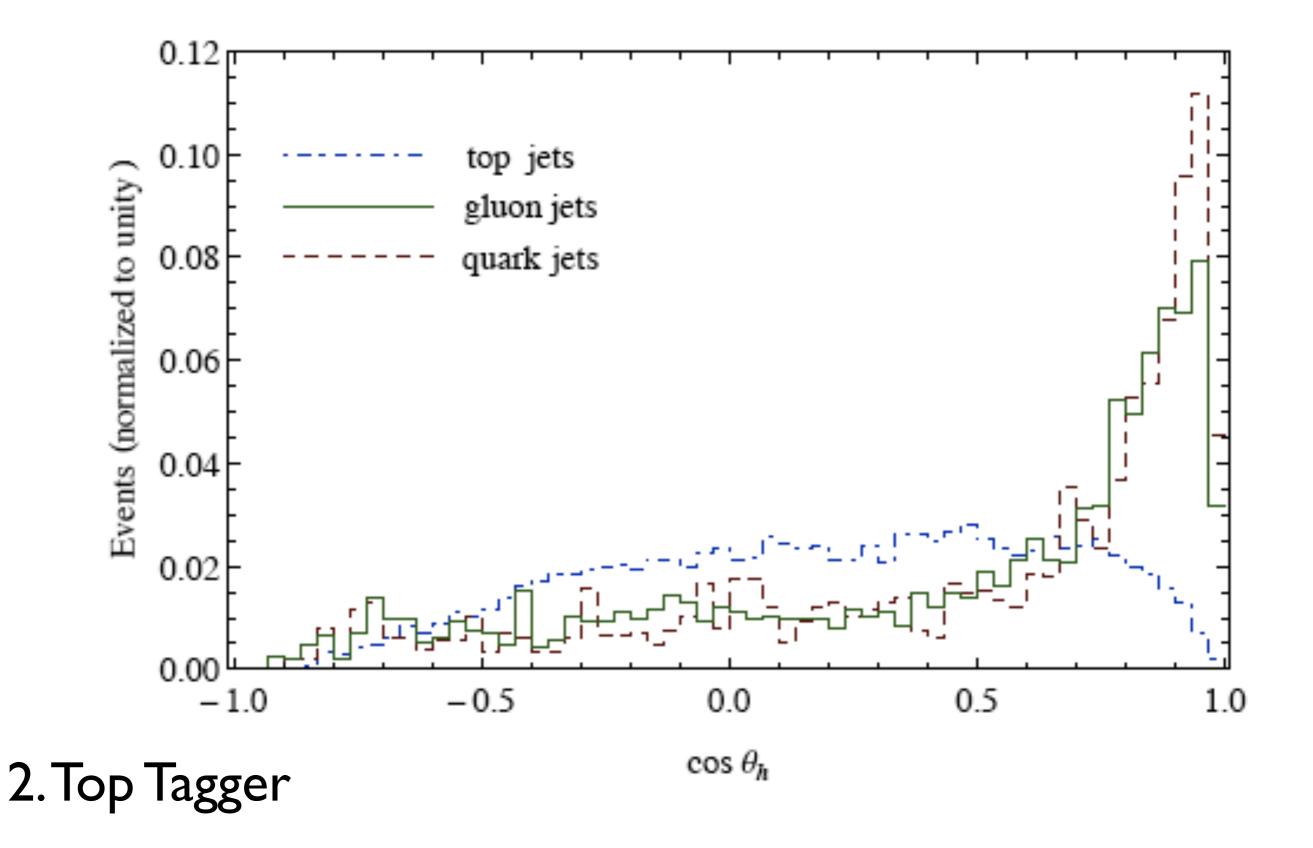
Top quarks can be involved into most new physics processes.

Hadronic Top tagger: 0806.0848, D.E. Kaplan, K. Rehermann, M.D. Schwartz, B. Tweedie



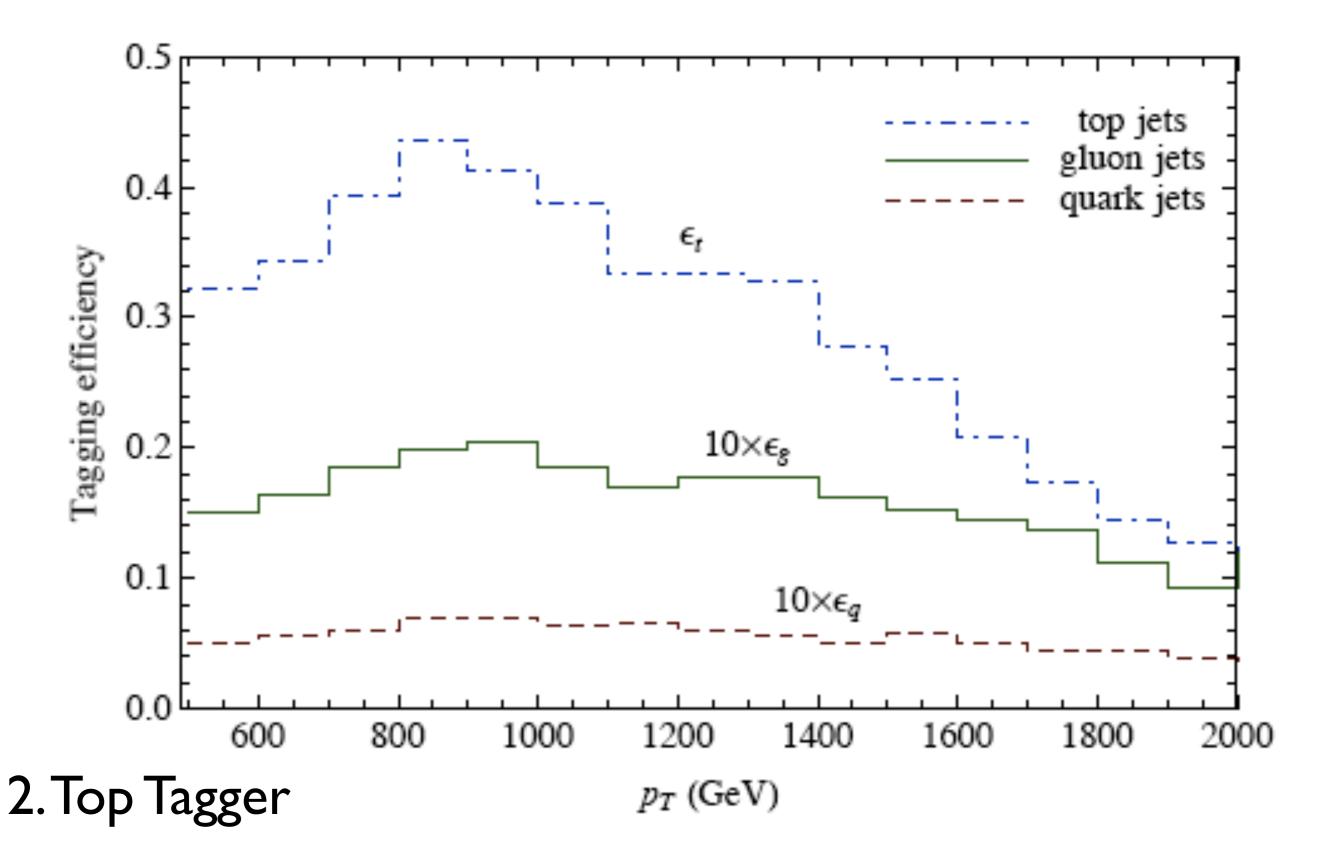
A hadronic top tagger exploits masses of top and W boson and helicity of W inside the top jet

Hadronic w helicty

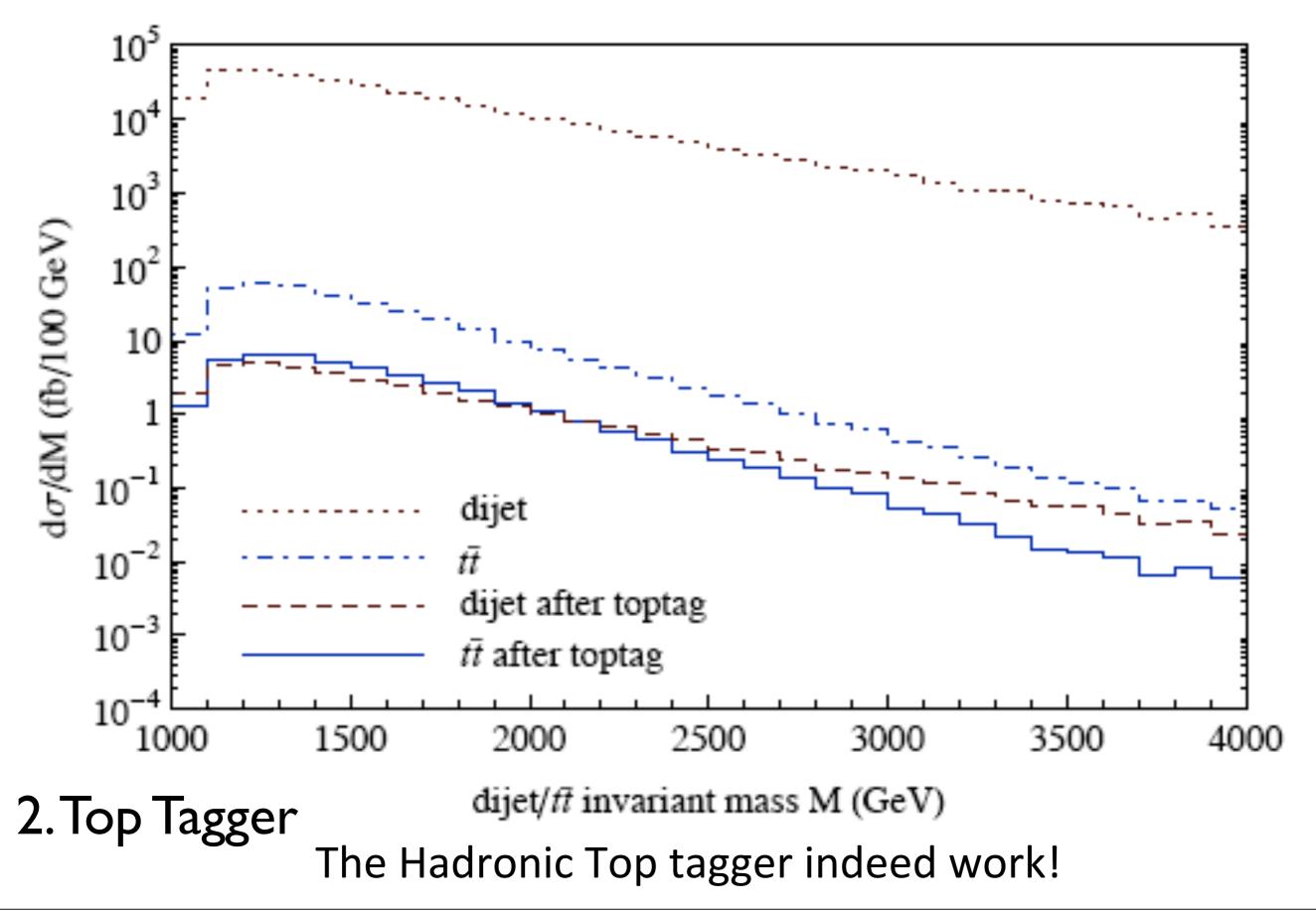


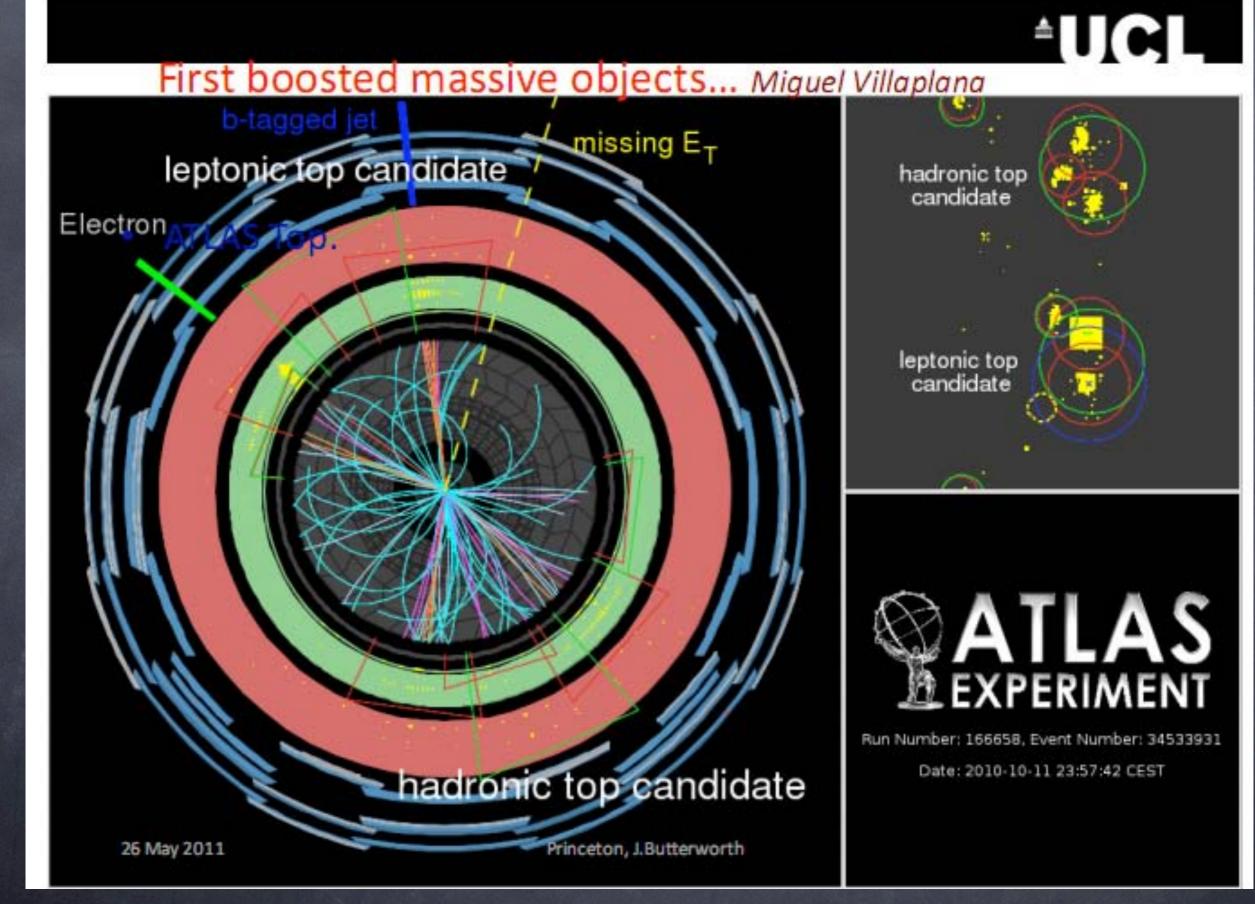
Top tagger: 0806.0848 D.E.Kaplan,K.Rehermann,M.D.Schwartz,B.Tweedie

Mt, mw, Hw, tagging efficiency



Hadronic Top tagger: 0806.0848 D.E.Kaplan,K.Rehermann,M.D.Schwartz,B.Tweedie

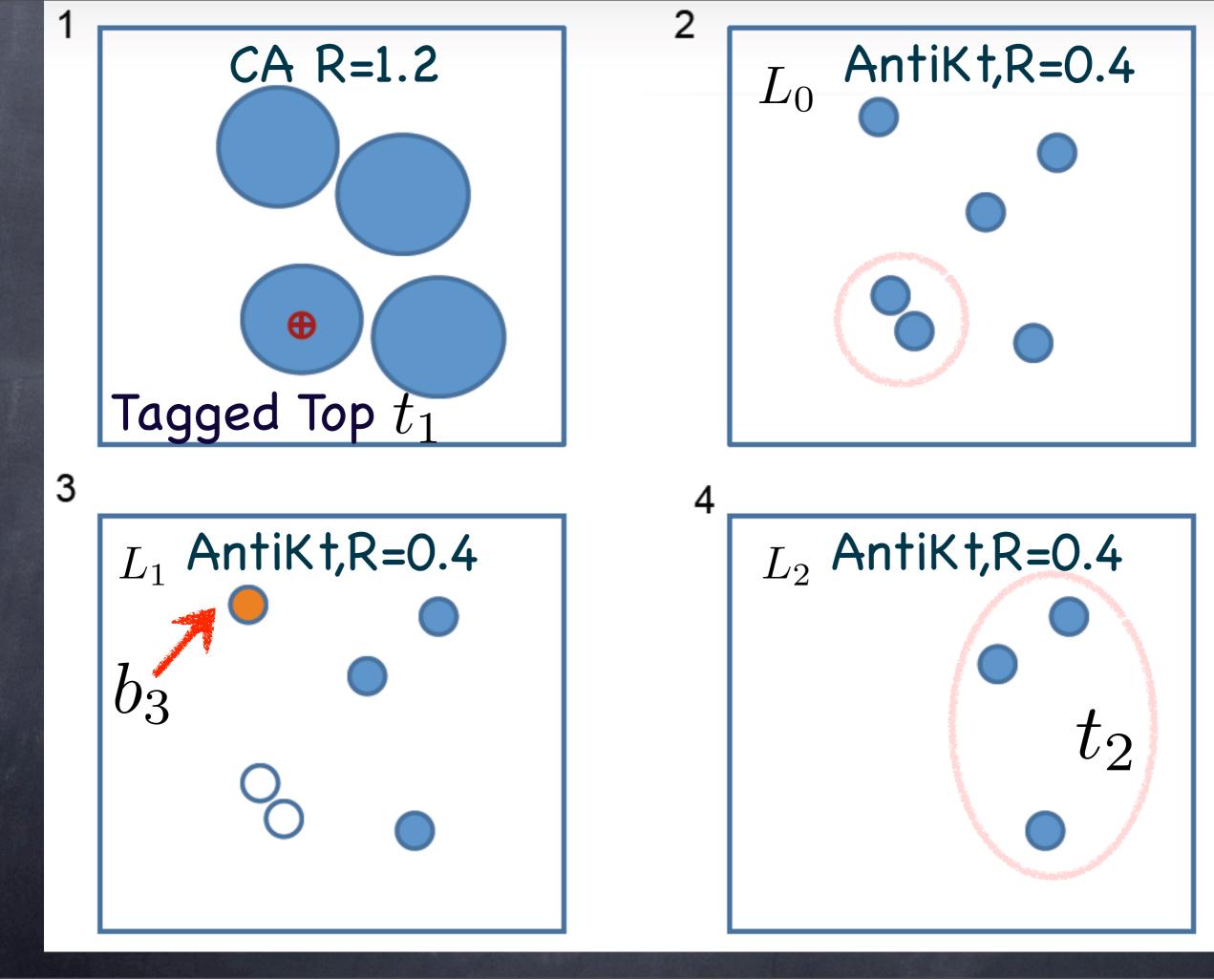




#### Boosted Top event observed at ATLAS

We propose a Hybrid-R method

- CA jet algorithm with a larger cone and top tagger to capture a highly boosted top  $t_1$ .
- Recluster the pseudo-jets in the event with a small cone with anti- $K_t R = 0.4$  and get a jet list  $L_0$ .
- If small-size jets are within the larger cone of the direction of the tagged top jet, remove them and get a new jet list  $L_1$ .
- Identify the most energetic jet in  $L_1$  as the  $b_3$  jet. The rest of unused jets form a jet list  $L_2$ .
- There are at least 3 jets in the  $L_2$ . Reconstruct the second top and W by  $\chi^2 = \frac{|m(j_1, j_2) - m_W^{PDG}|^2}{\sigma_W^2} + \frac{|m(j_1, j_2, j_3) - m_t^{PDG}|^2}{\sigma_t^2}$



R	0.7	0.8	0.9	1.0	1.1	1.2	1.3	1.4
$k_t$ algorithm	16%	22%	25%	28%	29%	29%	28%	27%
CA algorithm	13%	19%	24%	27%	29%	30%	30%	30%

#### $m_{H^+} = 1 \text{ TeV}$

R	0.7	0.8	0.9	1.0	1.1	1.2	1.3	1.4	
$k_t$ algorithm	25%	27%	28%	28%	27%	26%	24%	22%	
CA algorithm	23%	27%	29%	30%	30%	30%	30%	29%	
$1 \sqsubset \nabla_{2} V$									

#### $m_{H^+} = 1.5 \text{ TeV}$

Top tagger is better with larger mass when R fixed. CA algorithm is chosen for hadronic top tagger. Cone Size shall be optimized for S/B.

 $\mathcal{L} = \mathcal{L}_{SM} + \partial_{\mu}H^{+}\partial^{\mu}H^{-} - m_{H^{\pm}}^{2}H^{+}H^{-}$  $+H^{+}\bar{t}(Y_{L}P_{L} + Y_{R}P_{R})b + h.c. .$ 

 $Y_L = Y_R = 1$ 

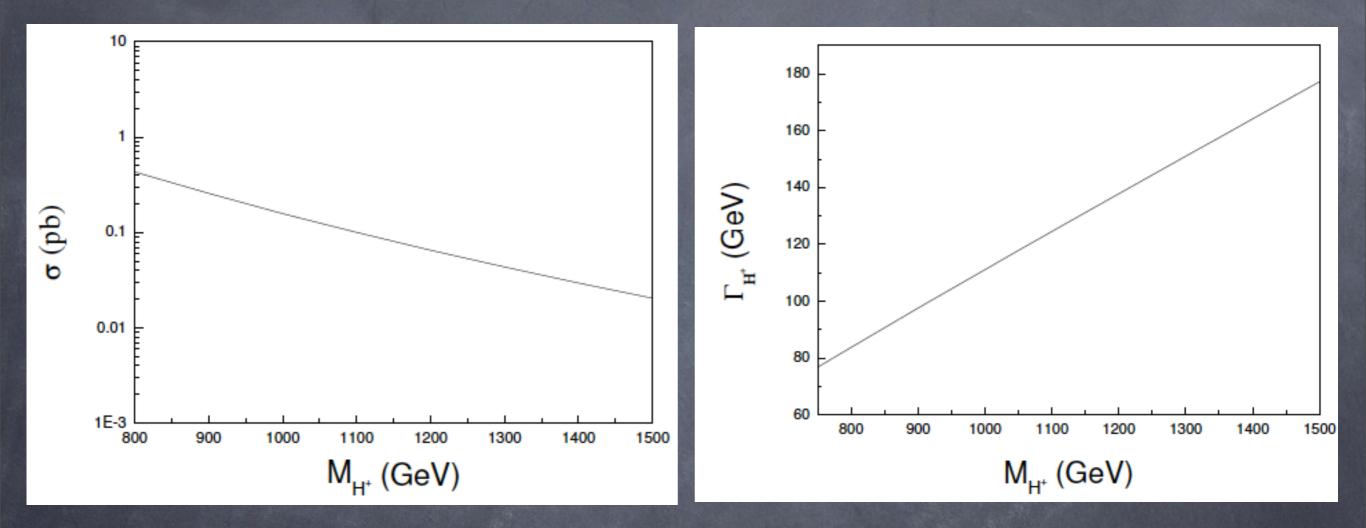
Must be tagged!

Conventions:

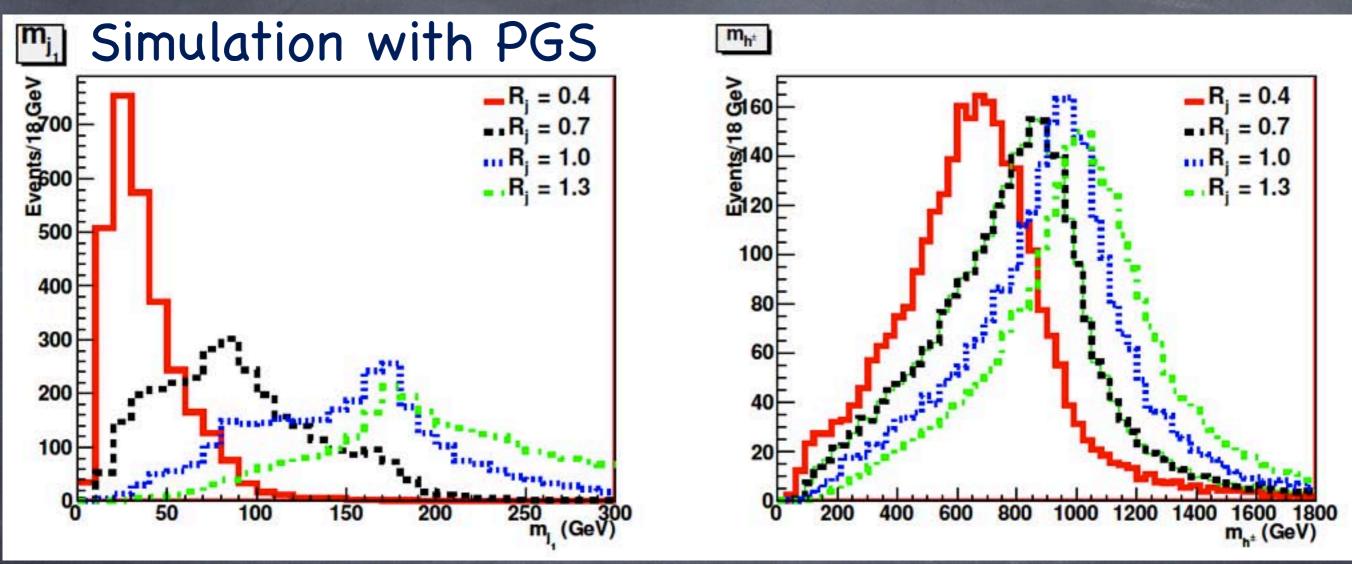
$$pp \to t_2 H^- \to t_2 t_1 b_3 / pp \to t_2 H^- b_4 \to t_2 t_1 b_3 b_4$$

# $pp \rightarrow Ht^-$ with s = 14 TeV K factor included

# The decay width $Br(H^+ \rightarrow tb) = 100\%$



The cross section is large enough, and the resonance is narrow enough to be resolved by detectors.



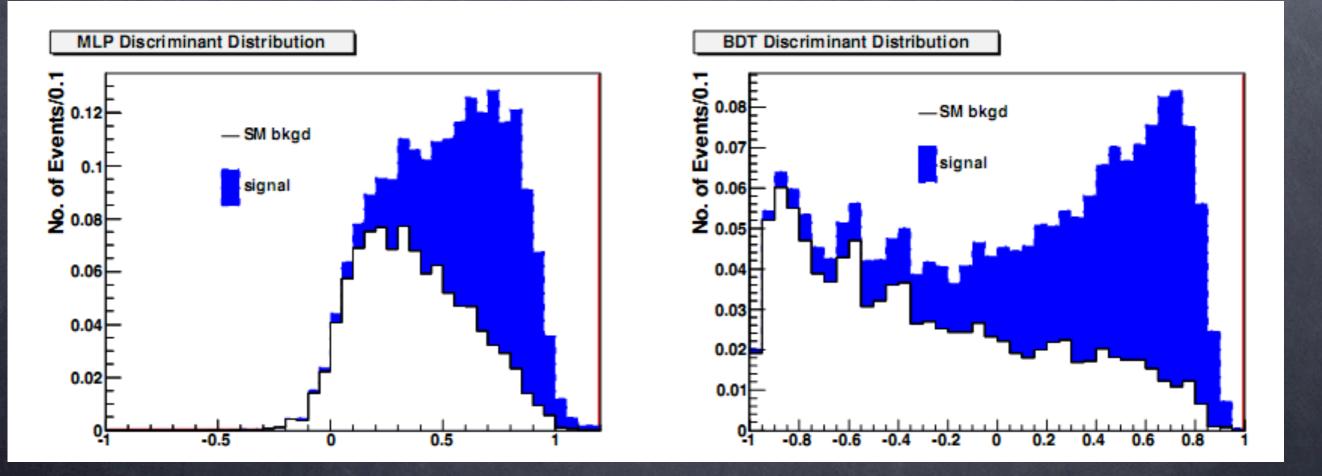
- Cone size and algorithm are crucial for the performance of the hadronic top tagger
- Charge Higgs mass bump can be approximately reconstructed.

• The detector effect and the decay width effect can not wash out signal

	signal	$t\bar{t} + jets$	QCD
	$m_{H^{\pm}} = 1 \text{ TeV}$		$n_j \ge 4$
Cross Section x Br (pb)	0.053	553	9186.8
After JH tagger & $H_T > 400 \text{GeV}$	29%	7%	$4.8 \times 10^{-3}$
Two <i>b</i> -taggings	10%	$4.2 \times 10^{-4}$	$4.8 \times 10^{-7}$
The number of jets in jet list $L_0, n_j \ge 6$	10%	$3.7 \times 10^{-4}$	$3.4 \times 10^{-7}$
$H_T > 700 \text{GeV} \& C > 0.3$	9.6%	$2.1 \times 10^{-4}$	$3.0 \times 10^{-7}$
The leading jet $E(j_1) > 350 \text{GeV}$	8.9%	$1.6 \times 10^{-4}$	$2.5 \times 10^{-7}$
The second leading jet $E(j_2) > 250 \text{GeV}$	7.9%	$1.3 \times 10^{-4}$	$2.2 \times 10^{-7}$
$P_t(b_3) > 300 \text{ GeV } \& P_t(t_1) > 300 \text{ GeV}$	5.3%	$4.5 \times 10^{-5}$	$7.7 \times 10^{-8}$
$ m_{W_2} - m_W^{\rm PDG}  < 20 { m ~GeV}$	3.1%	$2.6  imes 10^{-5}$	$2.3 \times 10^{-8}$
$ m_{t_2} - m_t^{\text{PDG}}  < 30 \text{ GeV}$			
$ m_{H^{\pm}} - m_{H^{\pm}}^{\text{assumed}}  < 200 \text{ GeV}$	2.5%	$1.5 \times 10^{-5}$	$1.5 \times 10^{-8}$
Events in 100 $fb^{-1}$	133	830	13.8

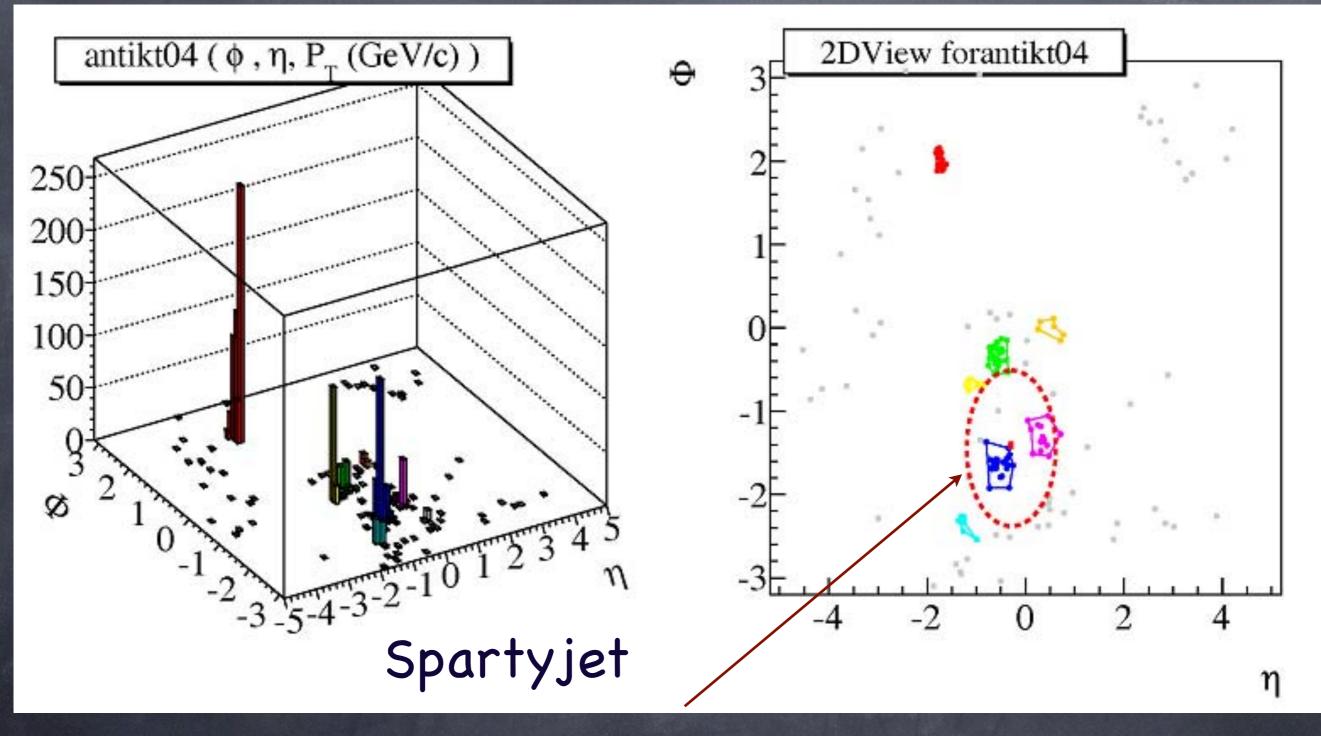
#### A Cuts based method can work but ...

	signal	$t\bar{t} + jets$	QCD
	$m_{H^{\pm}} = 1.0 \text{ TeV}$		$n_j \ge 4$
After simple cuts	2.5%	$1.5 \times 10^{-5}$	$1.5 \times 10^{-8}$
After NN cut $(NN > 0.6)$	5.5%	$2.0 \times 10^{-5}$	$3.0 \times 10^{-8}$
After BDT cut $(BDT > 0.5)$	5.7%	$2.1 \times 10^{-5}$	$3.2 \times 10^{-8}$



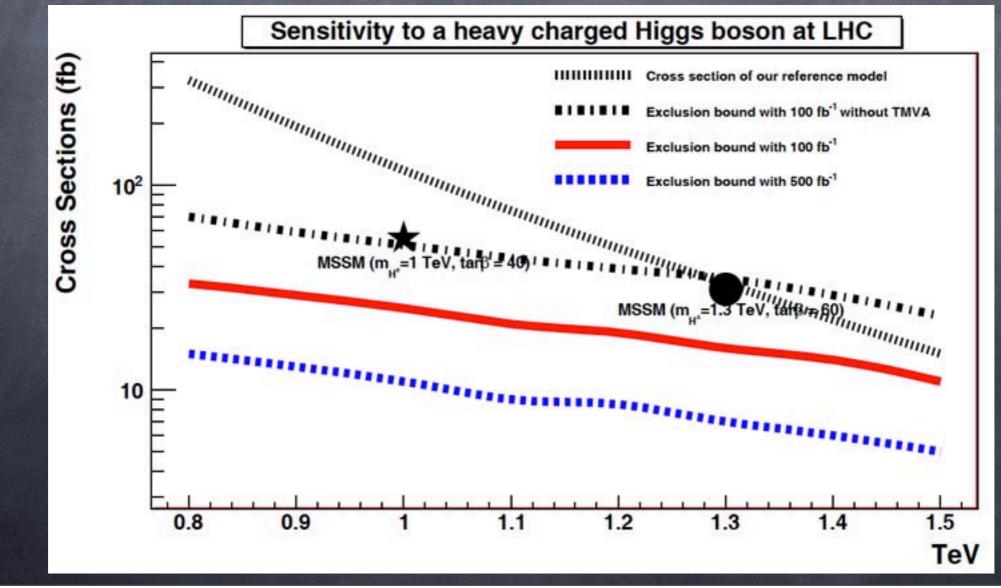
#### Significance can be enhanced 100% or so by TMVA.

#### A closer and intuitive look at a signal event.



#### Conventional top reconstruction method fails for tagged top

$m_{H^{\pm}}$ (TeV)	0.8	0.9	1	1.1	1.2	1.3	1.4	1.5
$\sigma$ (fb)		192	118	75	49	33	22	15
$\frac{S}{\sqrt{B}}$ (with two <i>b</i> taggings & TMVA)		13.3	9.5	7.0	5.1	4.0	3.1	2.6
lower bound on $\sigma(fb)$	33	29	25	21	19	16	14	11
$\frac{S}{\sqrt{B}}$ (with two <i>b</i> taggings without TMVA)		6.5	4.6	3.4	2.5	1.9	1.5	1.3
lower bound on $\sigma(fb)$		59	51	44	39	35	29	23



### 4. Conclusion

• We propose a hybrid-R method to detect heavy charged Higgs boson with hadronic top tagger(s).

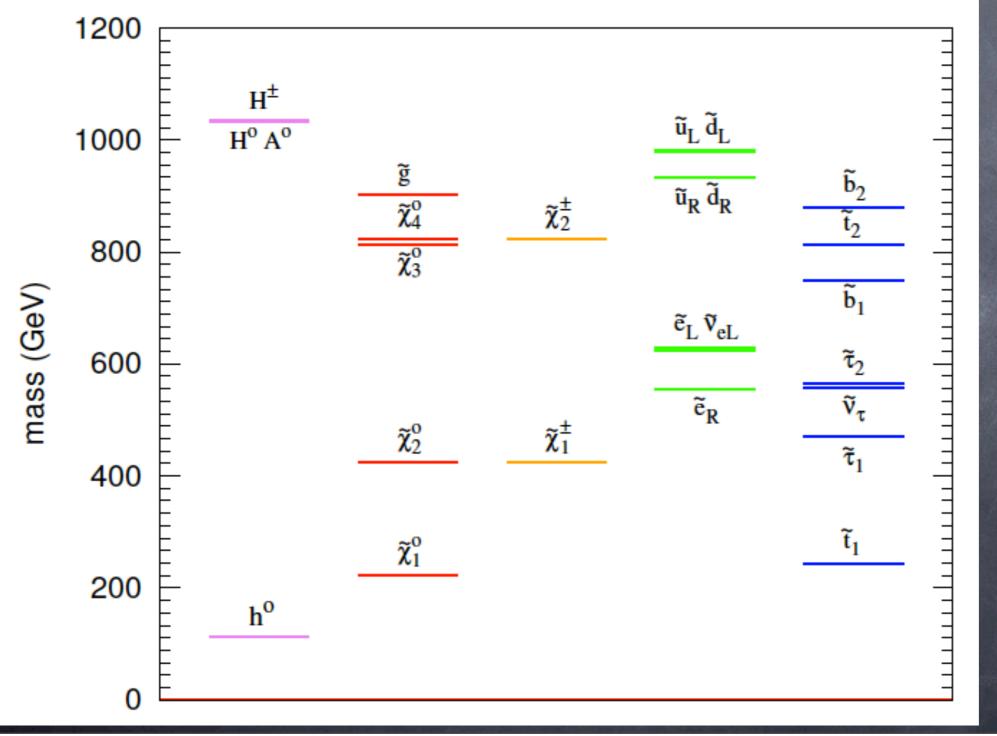
• The full hadonic mode is used as a test field and the method does work.

• The heavier the charged Higgs boson, the better performance is our method.

The heavy charged Higgs (from 0.8 TeV to 1.5 TeV) can be covered with 14 TeV and a 100 fb<sup>-1</sup> dataset just by using hadronic mode!

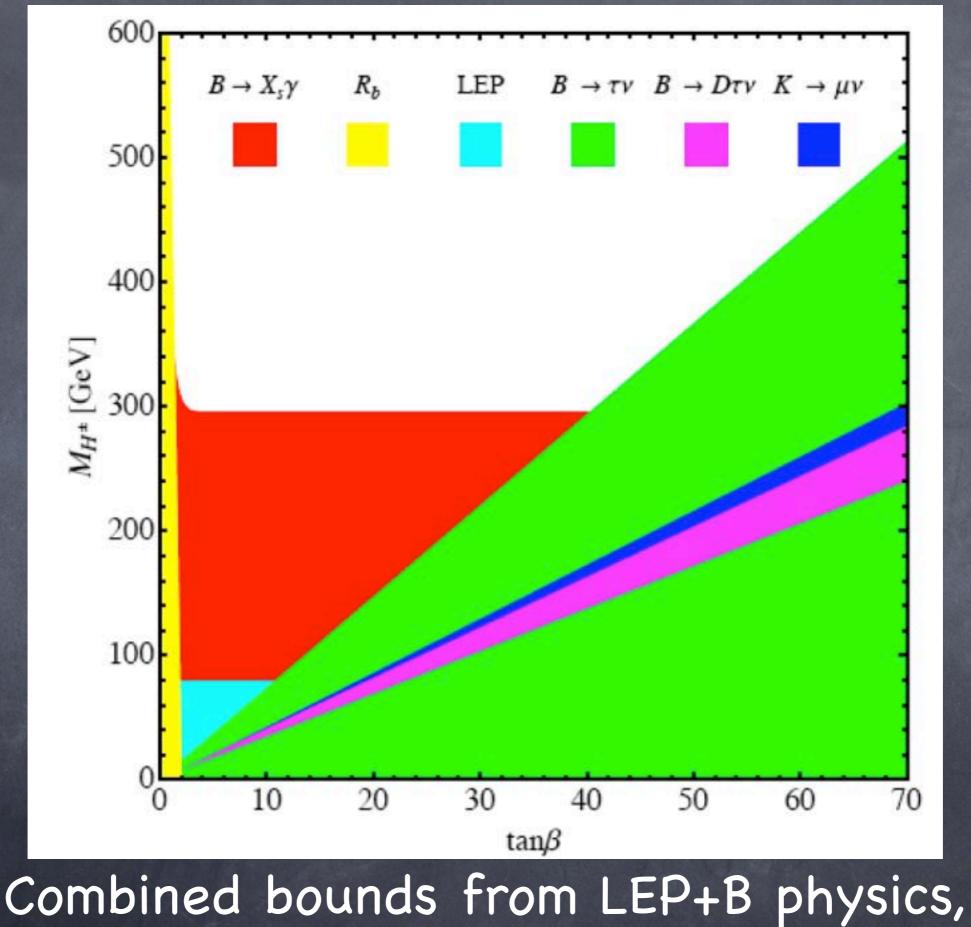
# Thanks you!

# 1. Why A Heavy Charged Higgs Boson?

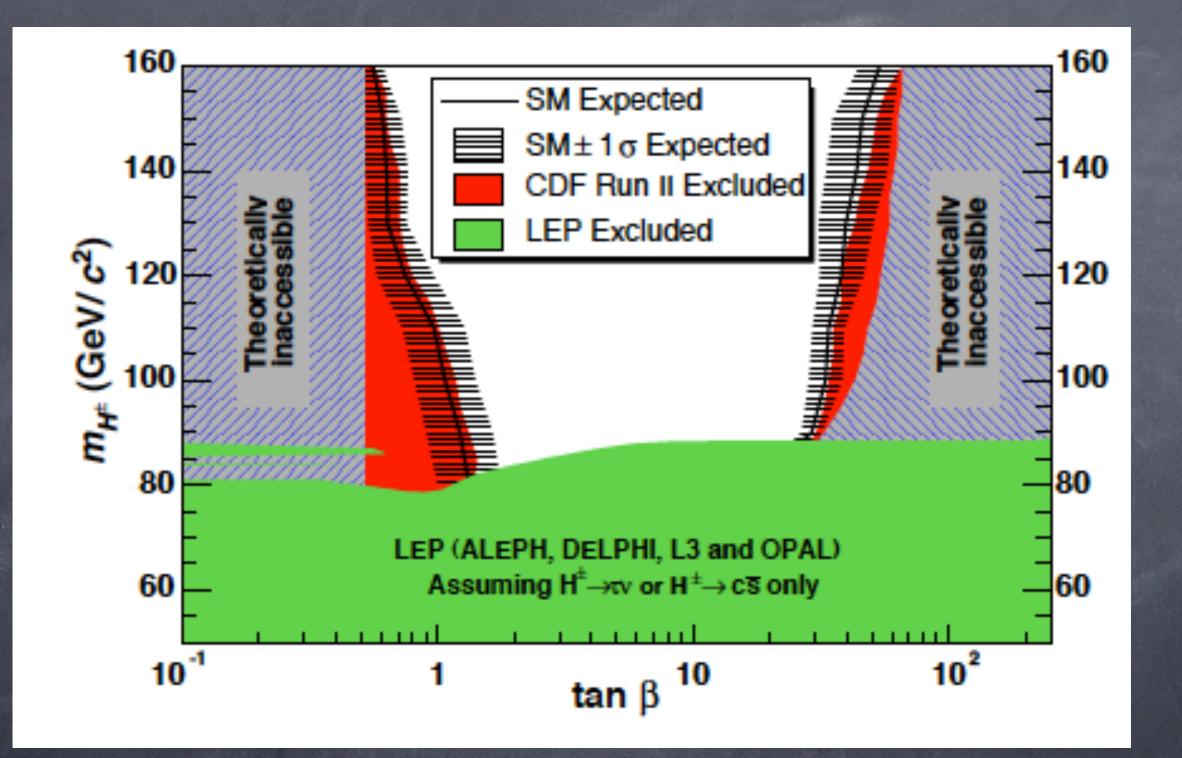


An example in a nonuniversal MSSM X.J.Bi,Q.S.Yan,P.F.Yin 1111.2250

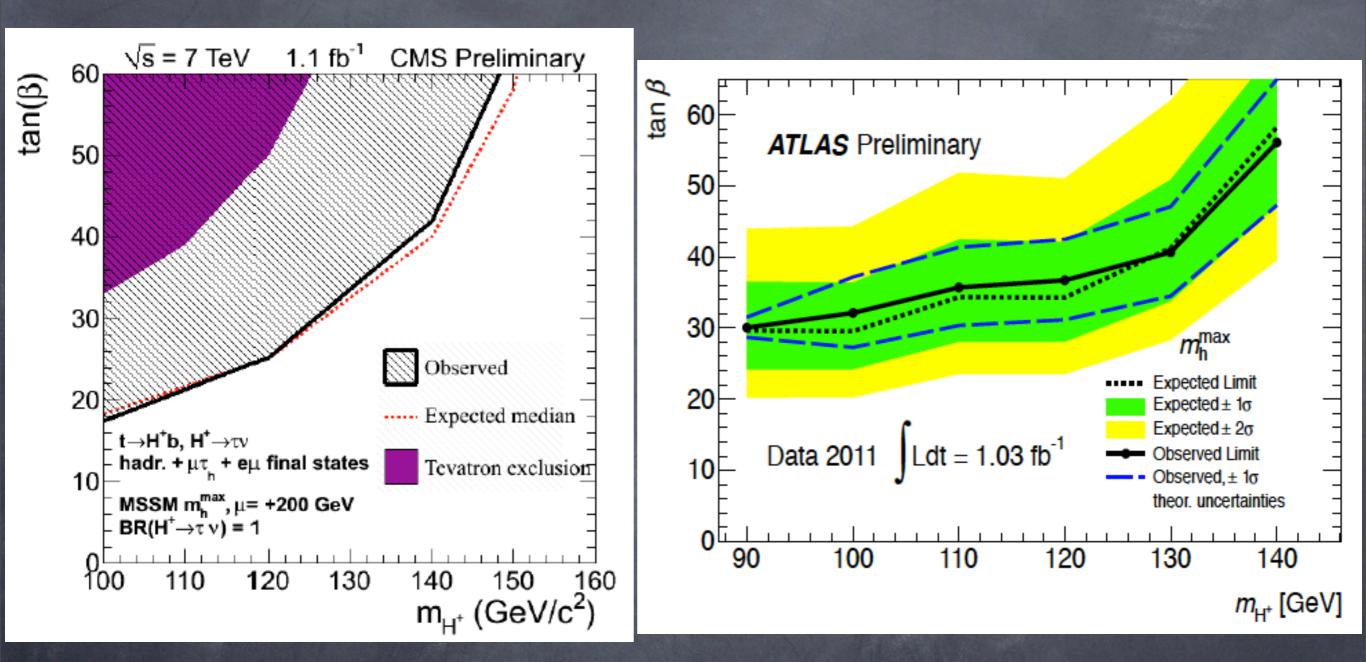
BMP 2



bench mark model (2HDM),0805.2141,Haisch

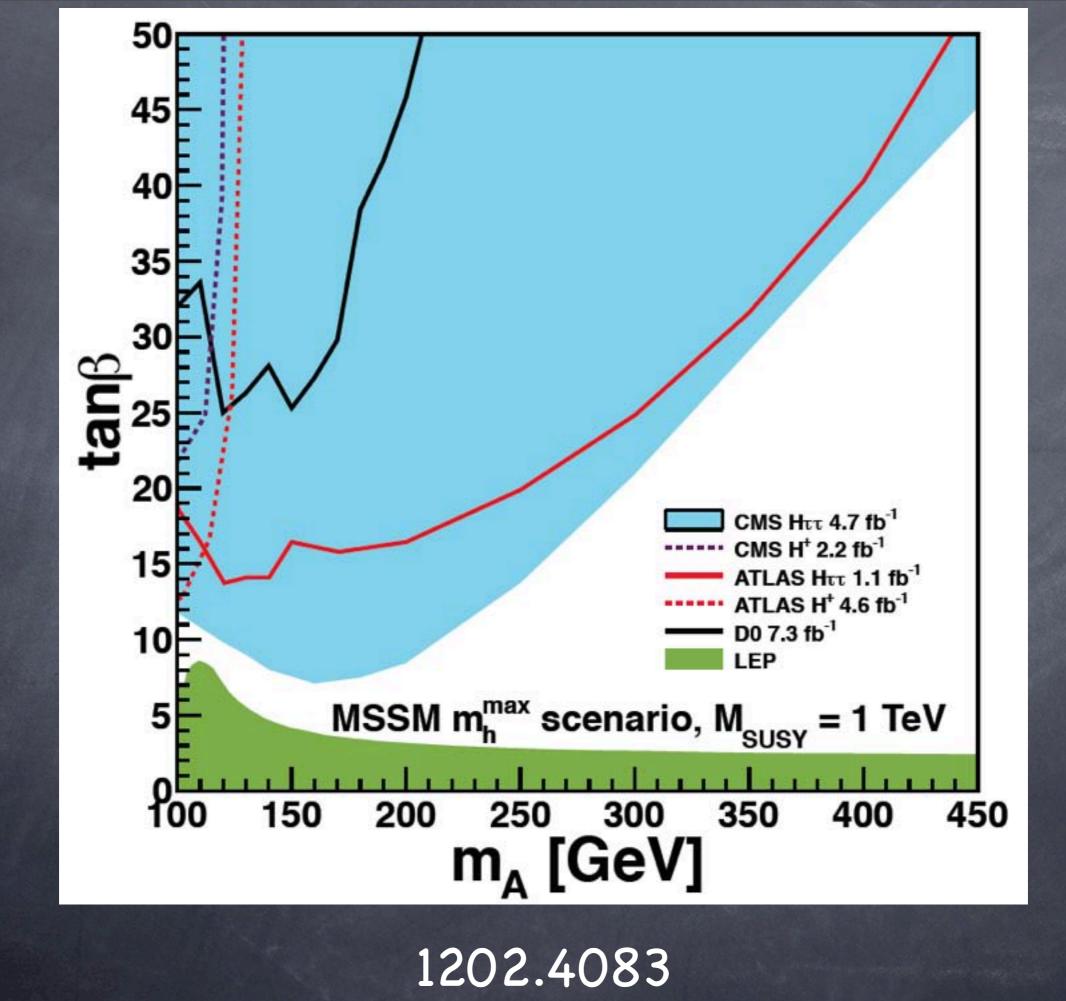


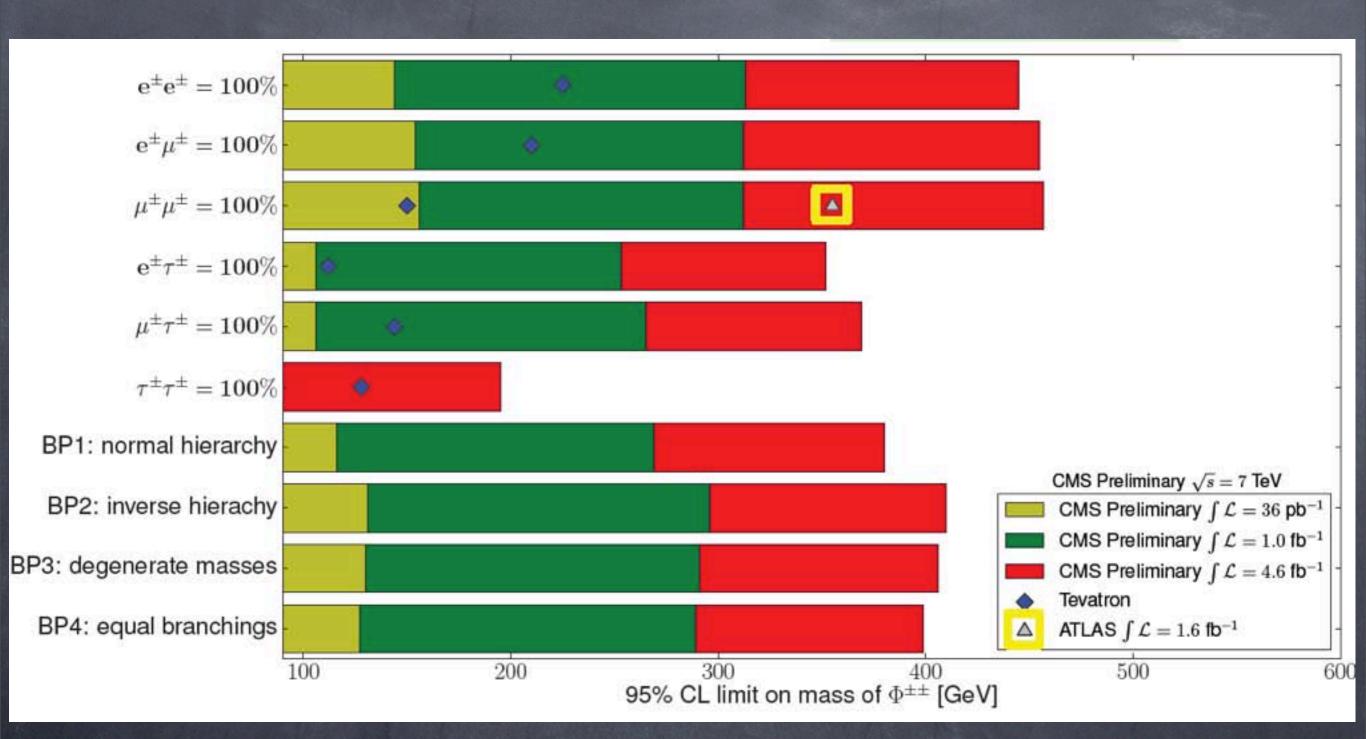
Bound to a charged Higgs boson(2HDM) from direct search from Tevatron&LEP



ATLAS,1201.5886

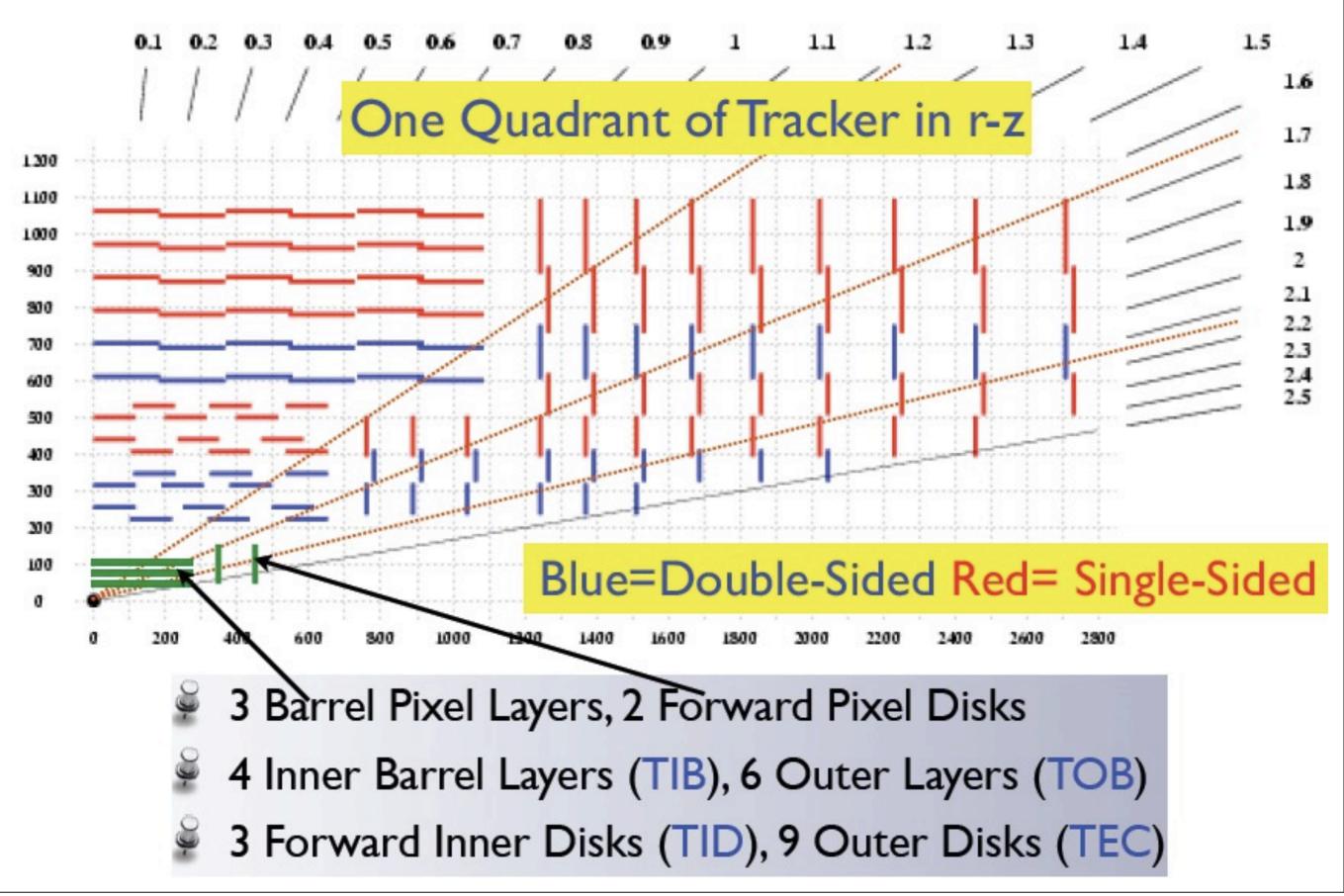
### CMS, 1201.4983

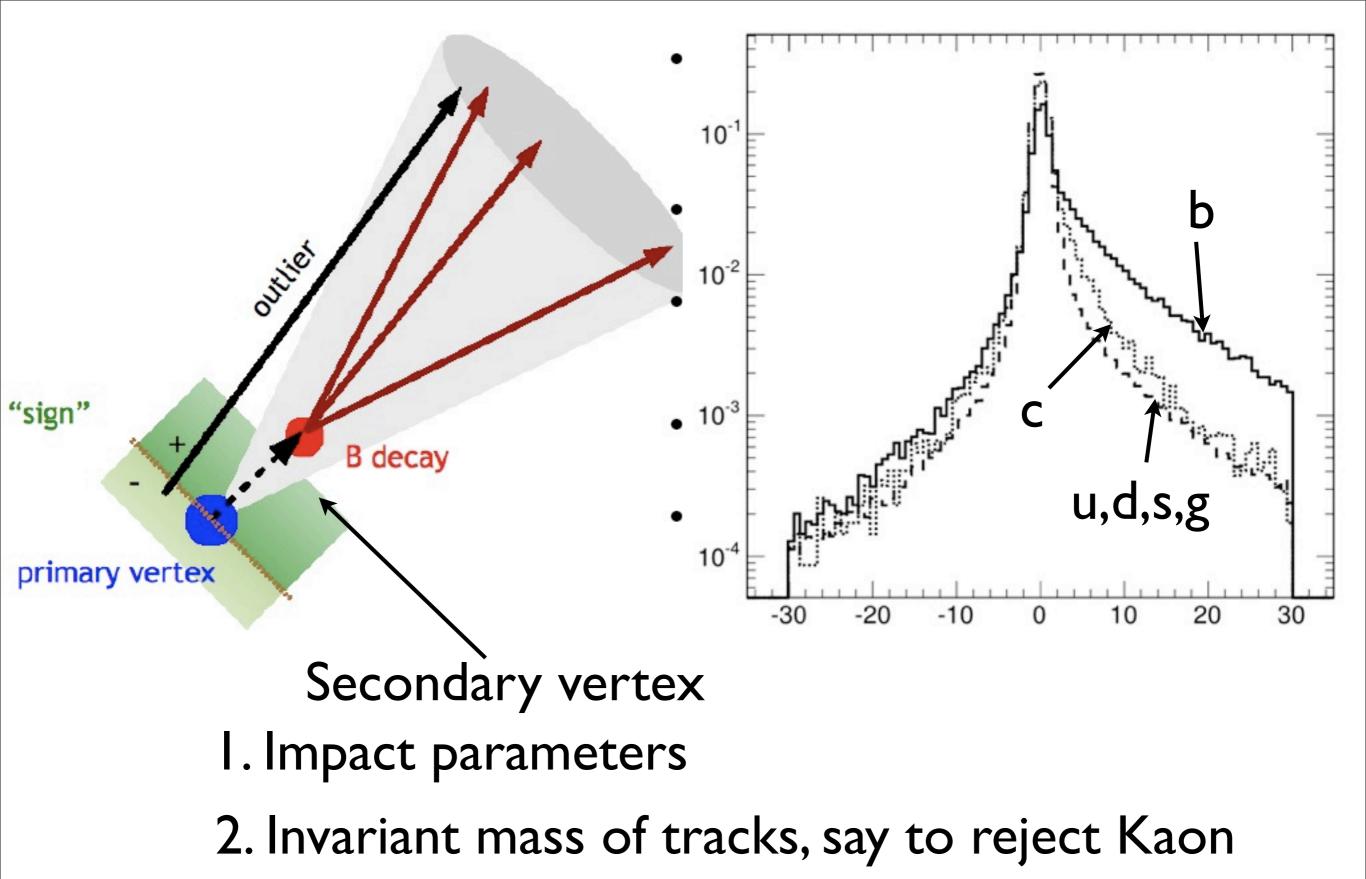




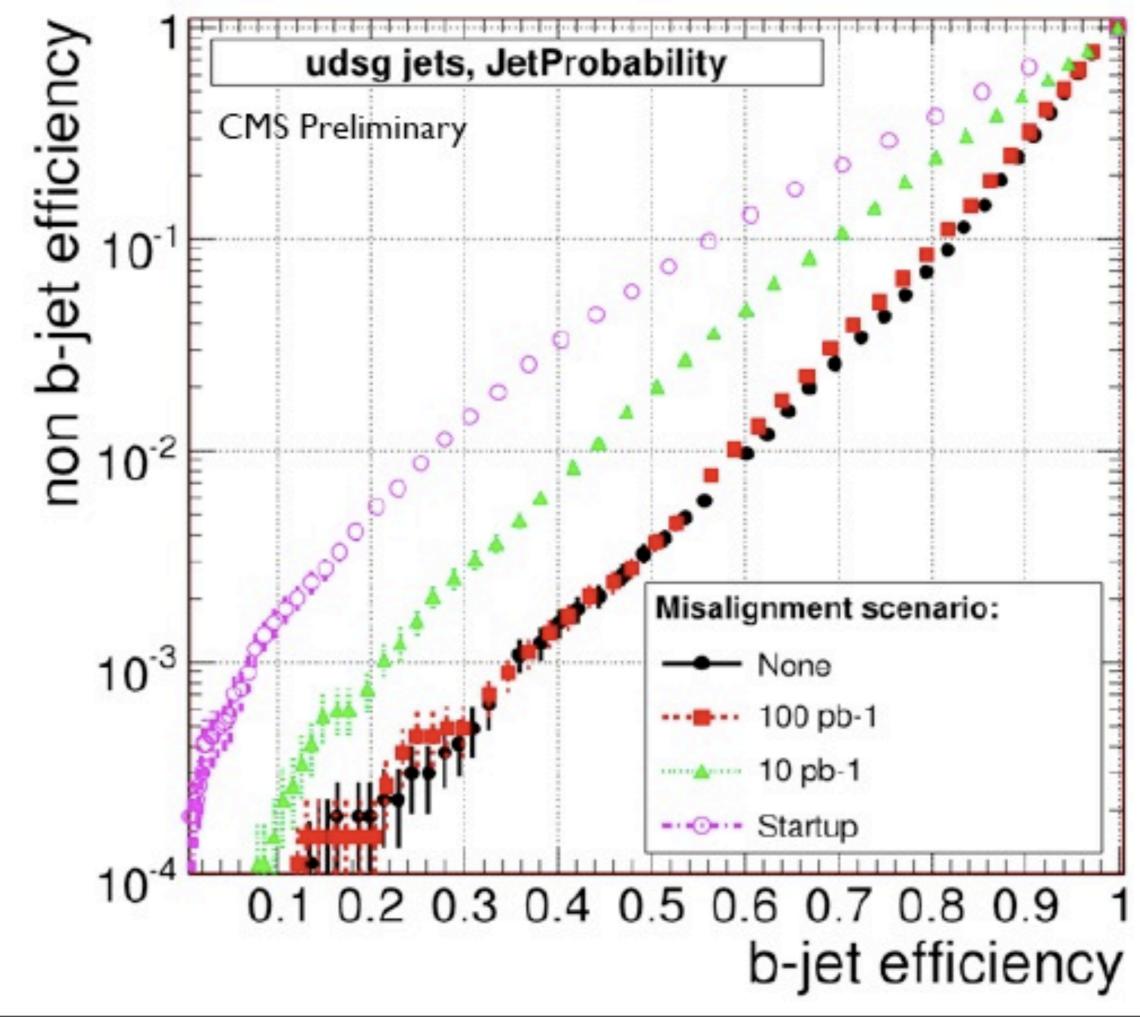
### Bounds for doubly charged Higgs boson

The layout of Si tracker @ CMS  $\delta L \sim \mu m$ 

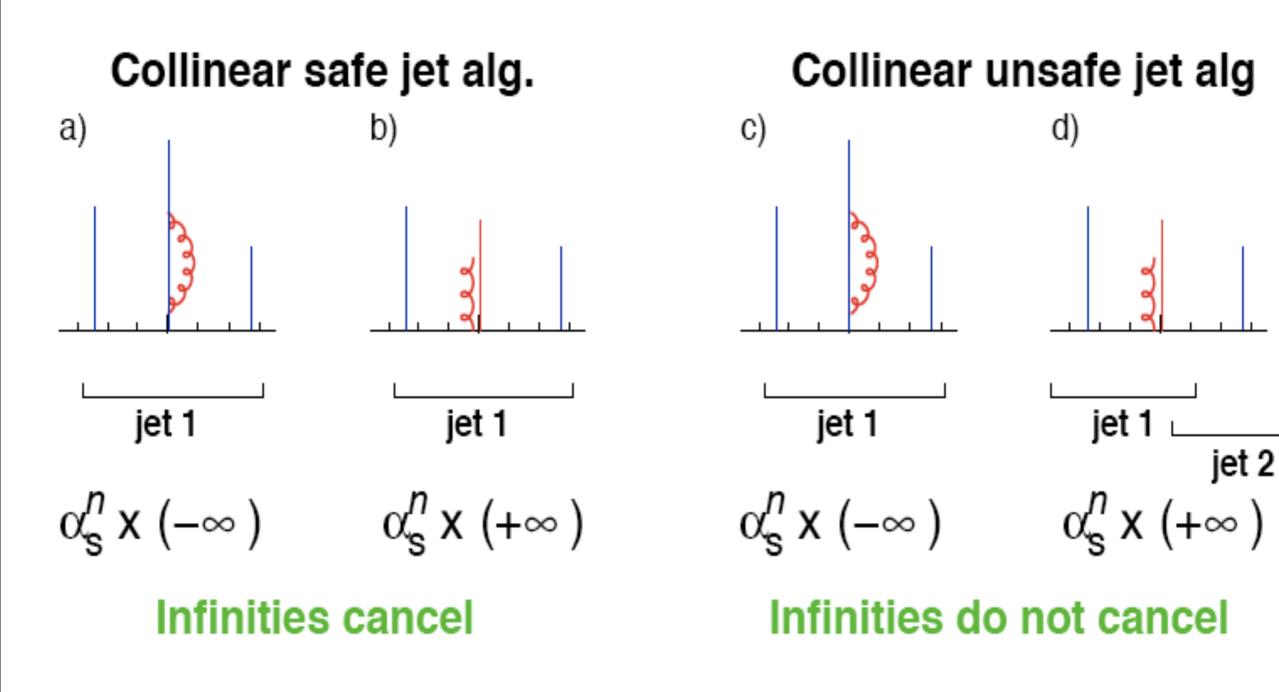




3. Charge multiplicity, b jet is massive, more charged particles can be produced.



Towards Jetography: 0906.1833, G.P. Salam



Towards Jetography: 0906.1833, G.P. Salam

<sup>A</sup> Cone algorithms

Sternman, Weinberg, 1977

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2 < R^2$$

$$\frac{\sigma_2}{\sigma} = 1 - \frac{32}{3} \frac{\alpha_s}{2\pi} \log \frac{1}{\delta} \log \frac{1}{\epsilon}$$

Intuitive, but not infrared safe.

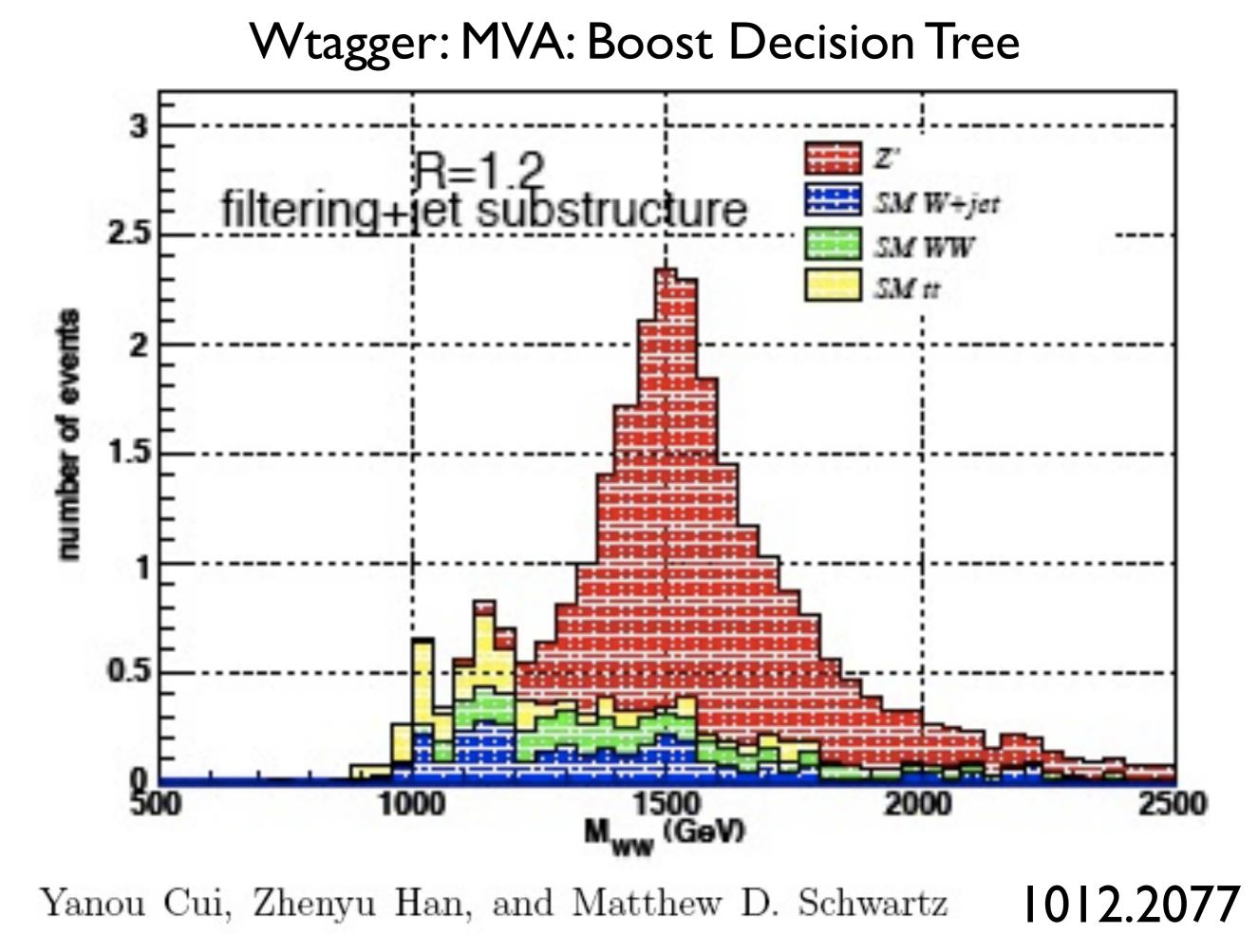
### <sup>B</sup> Sequential recombination jet algorithms

Jade 
$$y_{ij} = \frac{2E_iE_j(1-\cos\theta_{ij})}{Q^2}$$

$$\kappa_{t} \qquad y_{ij} = \frac{2\min(E_i^2, E_j^2)(1-\cos\theta_{ij})}{Q^2}$$

$$c/A \qquad v_{ij} = 2(1-\cos\theta_{ij})$$
Anti-K\_t 
$$d_{ij} = \min(p_{ti}^{2p}, p_{tj}^{2p})\frac{\Delta R_{ij}^2}{R^2}$$

P=-1,anti-Kt



	$p_T \simeq 500 { m GeV}$	$p_T \simeq 1000 { m ~GeV}$	$p_T \simeq 1500 { m ~GeV}$
W	76%	77	72
h	59	61	62

## W-Higgs tagger rate

	$p_T \simeq 500 \mathrm{GeV}$	$p_T \simeq 1000 { m ~GeV}$	$p_T \simeq 1500 \mathrm{GeV}$		
$\operatorname{quark}\to W$	6.5%	6.5	5.9		
quark $\rightarrow h$	6.8	5.6	5.8		
gluon $\rightarrow W$	10.4	8.3	7.4		
gluon $\rightarrow h$	10.5	8.8	7.4		
QCD mistag rate					

A.Katz, M. Son, B.Tweedie, 1010.5253

$$\left\langle m_{j}^{2}\right\rangle \simeq C_{i}\,\alpha_{s}\,p_{T,j}^{2}\,\Delta R_{j_{1}j_{2}}^{2}$$

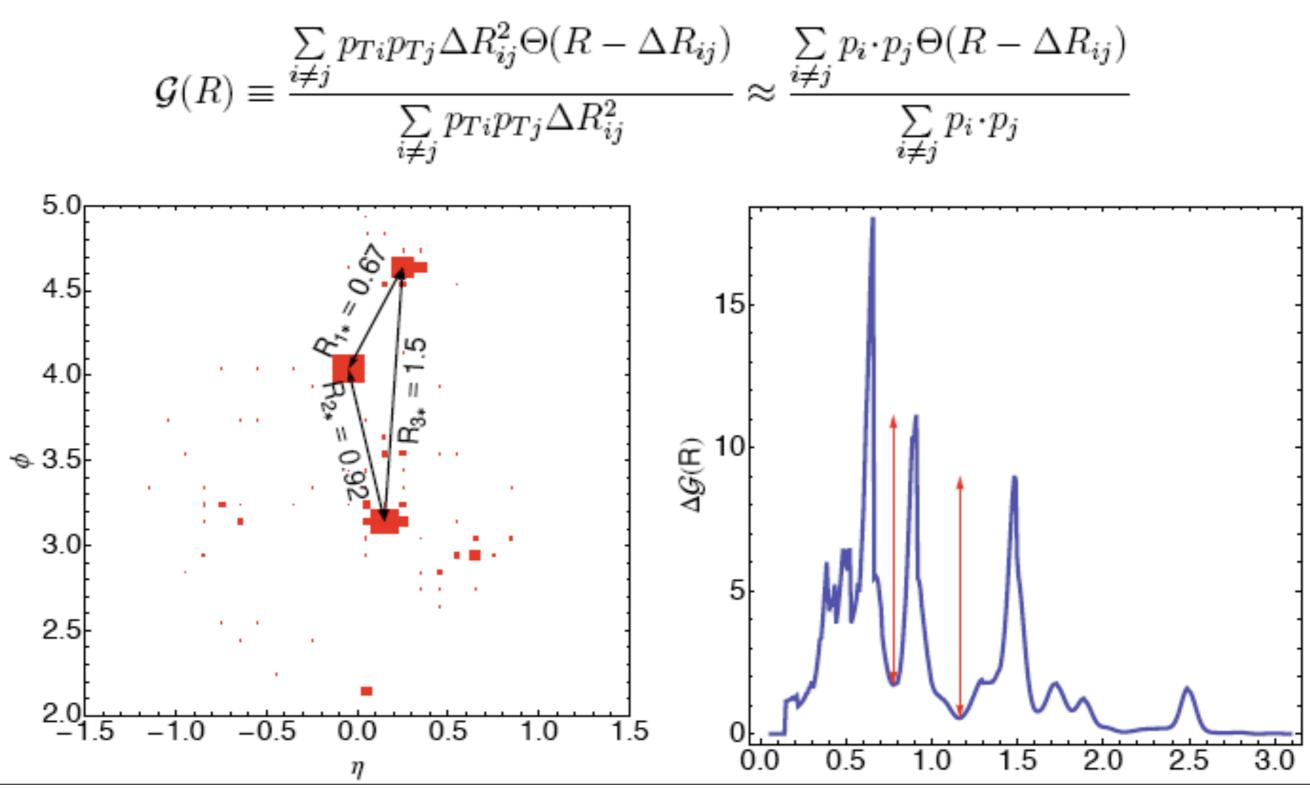
Gluon jet can be more massive than quark jet

S.D.Ellis, J.Huston, K.Hatakeyama, P.Loch, M.Tonnesmann, Prog. Part. Nucl. Phys. 60, 484

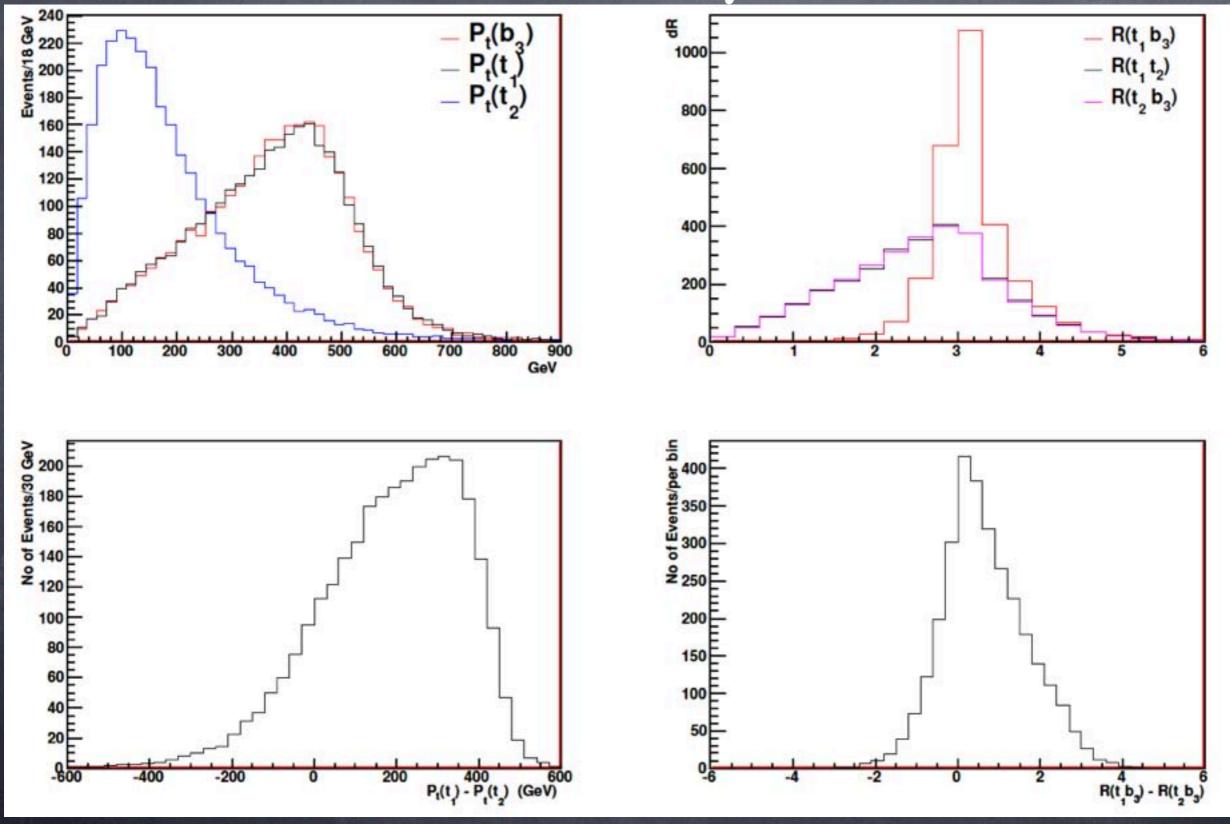
$$\left\langle \delta m_j^2 \right\rangle \simeq \Lambda_{\rm UE} \ p_{T,j} \ \left( \frac{R^4}{4} + \frac{R^8}{4608} + \mathcal{O}(R^{12}) \right)$$

The larger the cone size, the larger the contribution of underlying event. Pruning should be introduced to remove it.

M. Dasgupta, L.Magnea, G.P. Salam, JHEP0802,055

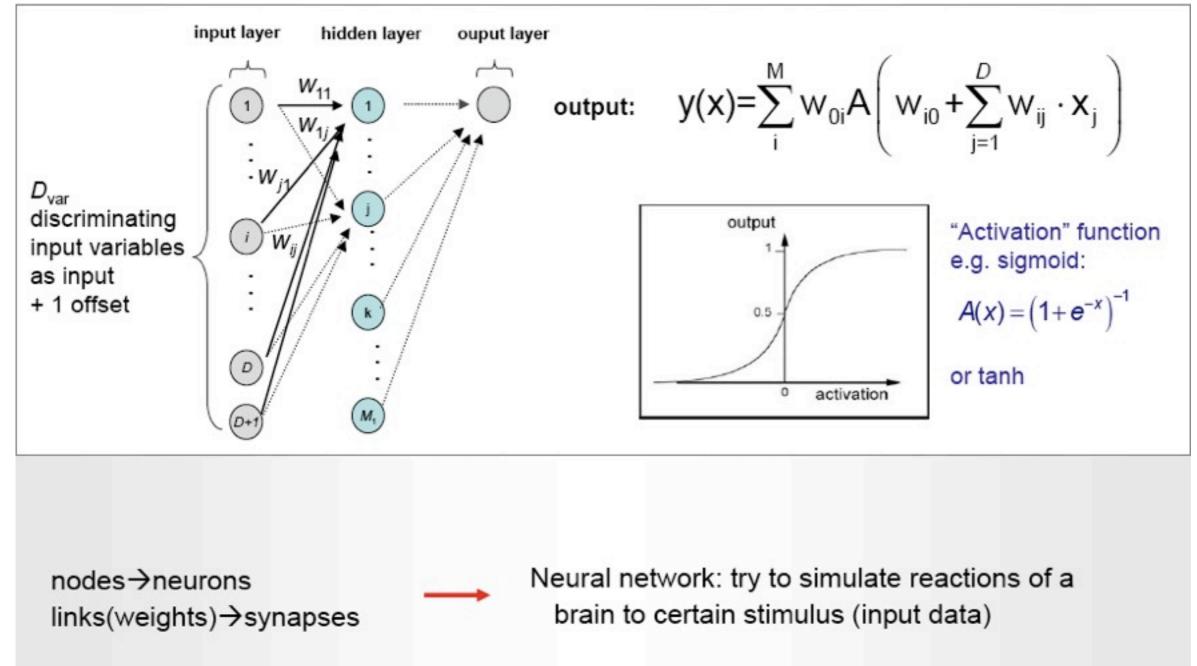


# 3. Feasibility@LHC

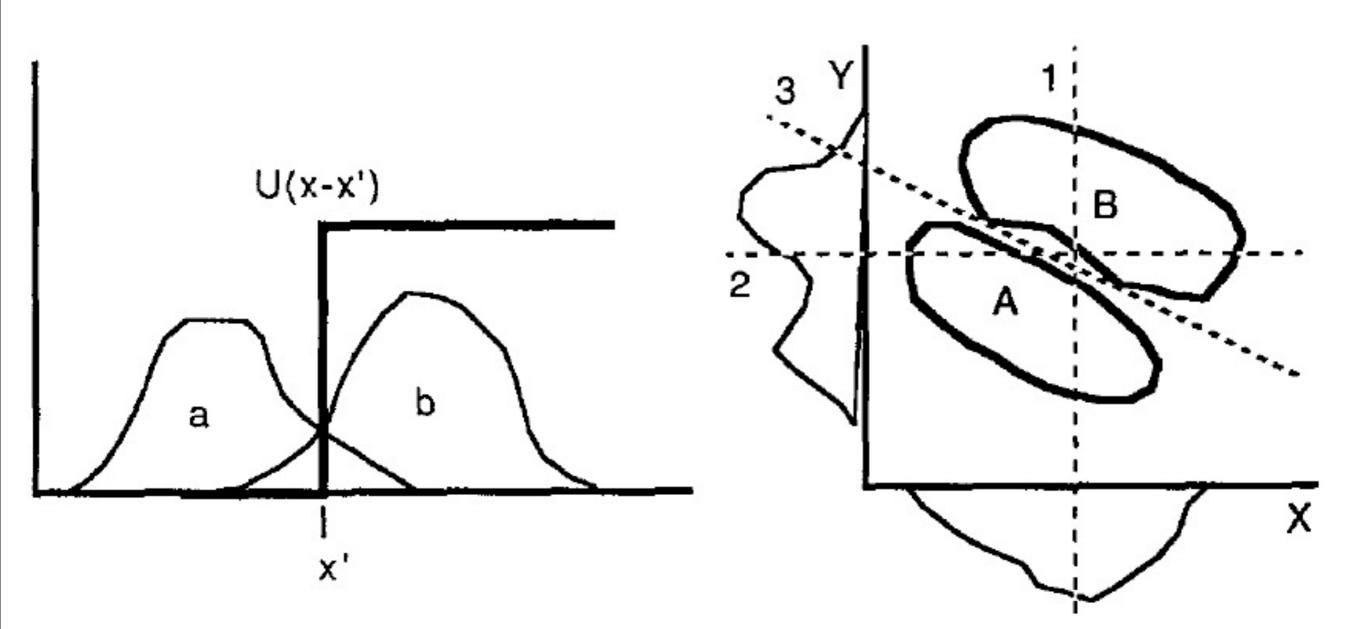


A Show Case  $m_{H^+} = 1$  TeV why we use those cuts?

# Neural Network Discrimination

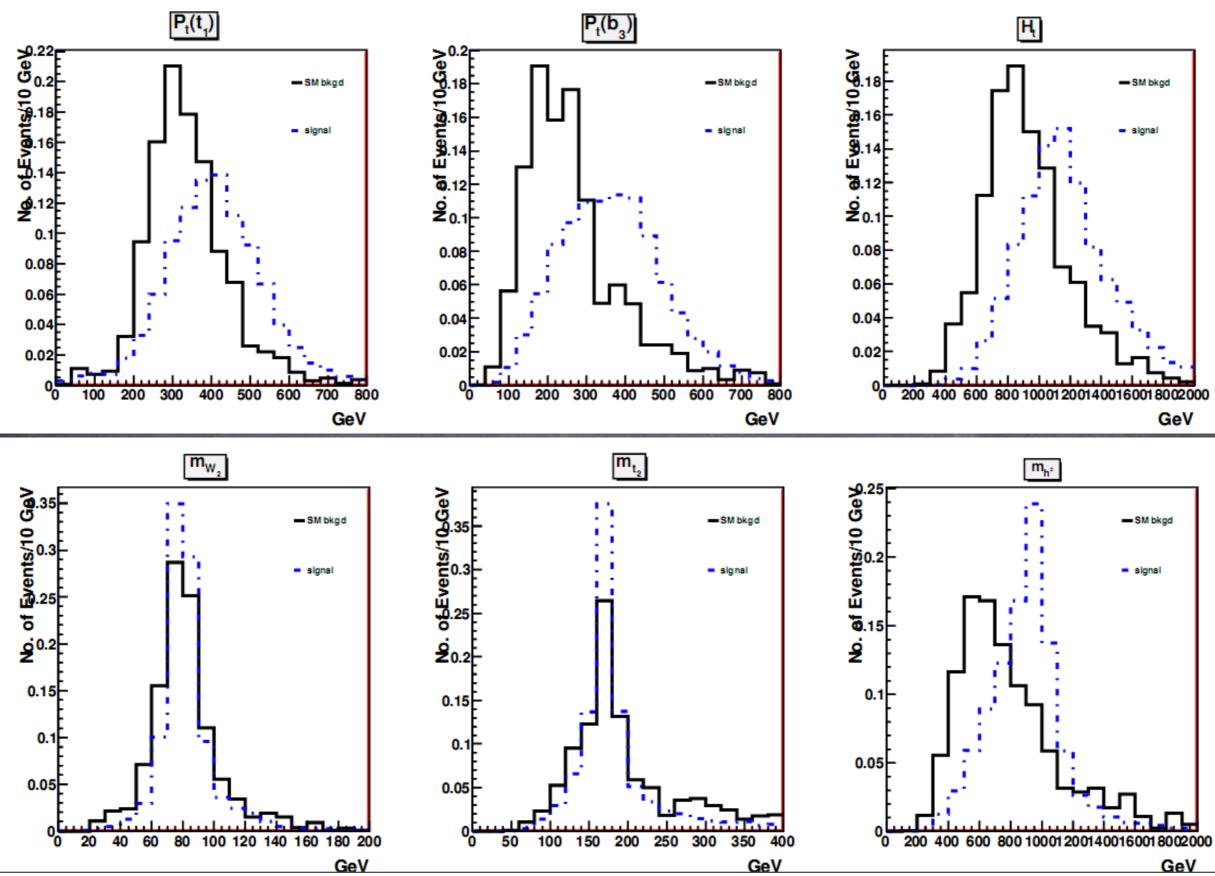


Considering the small number of signal compared with huge background we resort to the heavy machine gun: MVA



Why MultiVariate Analysis? To improve sensitivity of analysis when luminosity is limited and to reduce luminosity when budget fixed.

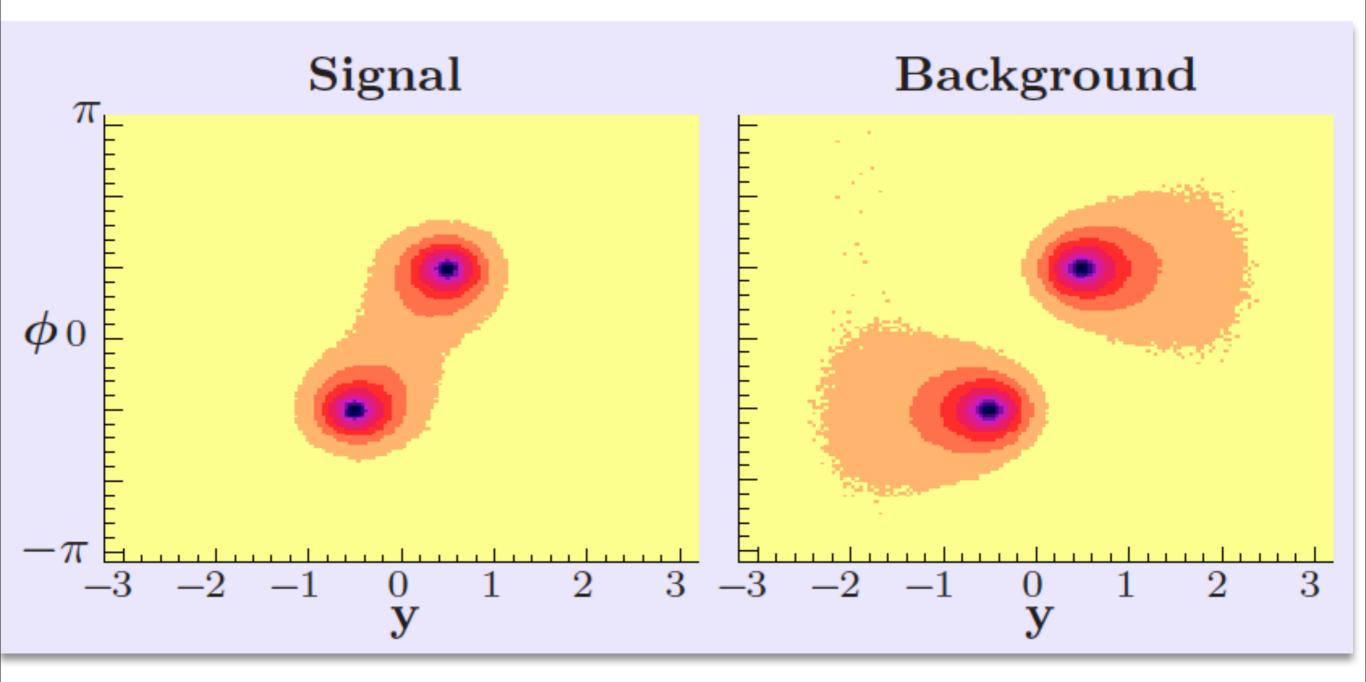
# 3. Feasibility@LHC



### **Color-Flow**

### Color Singlet

#### Color octet



J. Gallicchio, M. Schwartz,1001.5027 M.Janowiak,A.Hook,J.Wacker,1102.1012

$$\begin{split} \Psi(r) &= \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, r) \\ &\times \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \frac{d\sigma_{f}}{dP_{T}} \bar{J}_{f}^{E}(1, P_{T}, \nu_{\mathrm{fl}}^{2}, R, R) \right]^{-1} \\ & \left[ \sum_{f} \int \frac{dP_{T}}{P_{T}} \frac{d\sigma_{f}}{dP_{T}} \frac{d\sigma_{$$

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